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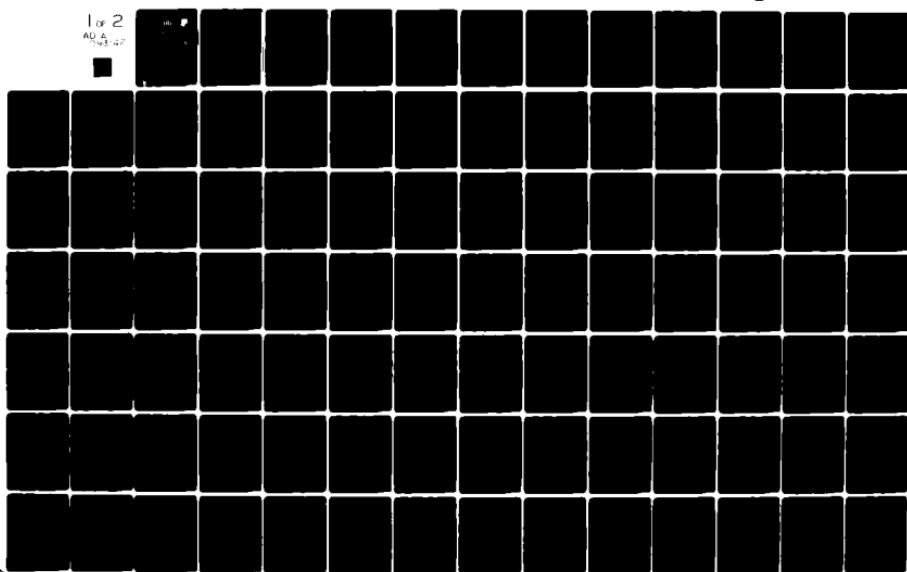
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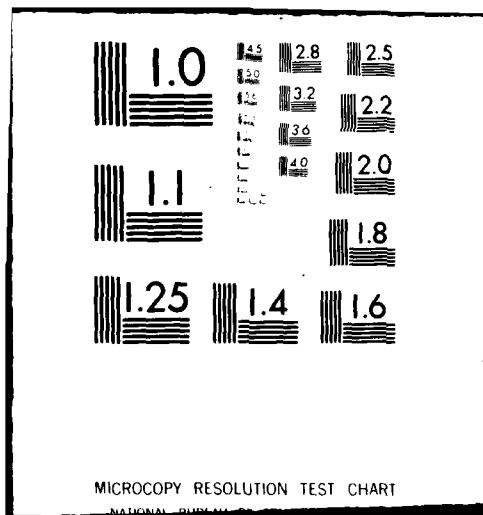
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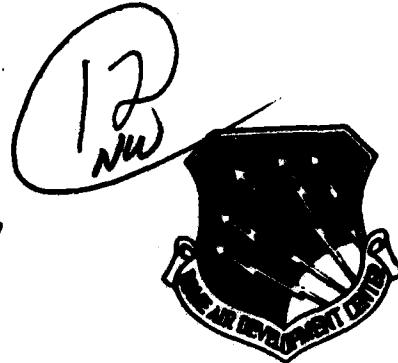




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RADC-TR-80-332
Final Technical Report
October 1980

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THEORETICAL ANALYSIS OF MULTIMODE FIBER STRUCTURES

EMTEC Engineering Incorporated

C. Yeh

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Some referenced figures do not appear in this document. The references are to computer calculations too voluminous to publish. The resulting analysis of the computer data is adequately presented in this report. Therefore, the missing data is considered irrelevant to the conclusions presented herein. The missing data may be obtained by contacting RADC (ESO) Hanscom AFB MA 01731.

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) This is a final report on the study of propagation characteristics of a light beam in multimode fiber structures. Realistic fiber structures made with commercially available fibers such as those provided by Corning or ITT were studied. The resultant computer programs may be used readily to generate design data for structures made with realistic fibers with step or parabolic index profiles. It is believed that our unique approach based on the scalar-wave FFT method may be extended to	20	

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L treat problems dealing with nonlinear fibers or fibers with frozen-in statistically varying index profiles.

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The work reported herein was supported by the Electronic Systems Division of the Air Force Systems Command, USAD, Hanscom AFB. The author wishes to express his special thanks to Dr. L. Eyes for his sincere interest in this project and for his suggestion in the application of our technique to study the large-size single mode fiber. Technical discussion with Dr. P. Gianino was also appreciated. Continuing support of this project by Dr. A Yang is gratefully acknowledged.

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E V A L U A T I O N

The effort as summarized in the report provides accurate theoretical analyses of the optical power transmission properties of a number of devices important to RADC efforts under TPO #4/D - Solid State Devices, Subthrust #3 - Electro-Optical Components. The devices include couplers, tapers, horns, and branches. The programs and techniques provided by the contractor permit in effect computer experiments to be done for a very large variety of design parameters. To do the actual experiments in the laboratory with this range of parameters would be enormously more expensive and time consuming. The report is then a crucial element in simplifying and accelerating the design process and in leading to final design specs in the shortest time.

Leonard Eyses
LEONARD J. EYGES
Project Engineer

I. INTRODUCTION

This final report summarizes the work performed under Contract F19628-80-C-0053 which the Electronic Systems Division of the Air Force Systems Command granted to the EMtec Engineering, Los Angeles, California. The work was begun in January, 1980 and completed in August, 1980.

The principal thrusts of this R & D study in performing numerical analysis of multimode fiber components were two fold. Firstly, we wish to learn the limitation (and possible improvement) of our numerical scheme¹ and secondly, we wish to obtain numerical data for realistic multimode fiber structures.

Specifically, the following tasks were carried out:

- a) Study the effect of step index gradient and of tight beam confinement by an adaptive coordinate scheme.
- b) Study the effect of the presence of absorber at the edge of the mesh on the beam propagation characteristics of multimode fiber structures.
- c) Compute the coupling characteristics of tapered multimode fiber couplers and unequal size fiber couplers .
- d) Obtain data for reflection coefficients and beam waist changes for multimode fiber tapers, horns and branches.

In section II we shall present the implementation of the adaptive coordinates in our numerical solution of the scalar wave equation. Then, the scheme to include an absorber at the edge of the mesh will be described. Finally, an approximate approach to obtain the reflection coefficients for complex fiber structures will be shown. Detailed results of our study on the proposed tasks are given in Section III. Concluding remarks and recommendations for future work are included in Section IV.

II. ANALYTICAL APPROACH

The basic approach taken to find the solution of wave propagation along complex fiber structures is to solve the reduced scalar wave equation via the fast Fourier transform (FFT) technique.² In this section we shall first indicate the conditions underwhich the exact vector wave equation may be simplified to yield the reduced scalar wave equation. Then we shall introduce the concept of adaptive coordinates³ and incorporate this concept in the solution of the reduced scalar wave equation via the fast Fourier transform technique.

A. Formulation of the Scalar Wave Approach. Starting with the vector wave equation for the electric field vector \underline{E} in the fiber structure,

$$\nabla \times \nabla \times \underline{E} - \omega^2 \mu_0 \epsilon \underline{E} = 0 \quad (1)$$

where ω is the frequency of the wave, μ_0 the permeability and $\epsilon = \epsilon(r)$, the inhomogeneous permittivity of the structure, and making use of the vector identity

$$\nabla \times \nabla \times \underline{E} = \nabla(\nabla \cdot \underline{E}) - \nabla^2 \underline{E} \quad (2)$$

and the relation

$$\nabla \cdot \underline{E} = - \frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E}, \quad (3)$$

one has

$$\nabla^2 \underline{E} + \omega^2 \mu_0 \epsilon \underline{E} - \nabla \left(\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) = 0 \quad (4)$$

Rewriting Eq. (4) gives

$$\nabla^2 \underline{E} + \omega^2 \mu_0 \epsilon_0 \left\{ \frac{\epsilon}{\epsilon_0} \underline{E} - \left[\frac{1}{\omega^2 \mu_0 \epsilon_0} \nabla \left(\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) \right] \right\} = 0$$

The relative importance of the terms within the curly brackets can be determined from the following

$$\frac{\epsilon}{\epsilon_0} \underline{E} = \mathcal{O}\left(\frac{\epsilon}{\epsilon_0} \underline{E}\right) \quad (6)$$

$$\frac{1}{\omega^2 \mu \epsilon_0} \nabla \left(\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) = \frac{1}{k_0^2} \mathcal{O}\left(\frac{\nabla \epsilon}{\epsilon} \cdot \nabla \underline{E}\right) = \mathcal{O}\left(\frac{\epsilon/\epsilon_0}{k_0 l} \underline{E}\right) \quad (7)$$

where the symbol \mathcal{O} means the "order of magnitude," and l is the smaller of the distance over which ϵ/ϵ_0 and \underline{E} change appreciably. For single-mode fiber structures, the values of ϵ/ϵ_0 and $k_0 l$ are typically in the range

$$\epsilon/\epsilon_0 = \mathcal{O}(2) \quad (8)$$

$$k_0 l = \frac{2\pi}{\lambda} l = \mathcal{O}(10^2 \text{ or } 10^3) , \quad l = \mathcal{O}(10\mu \text{ to } 100) \quad (9)$$

$$\lambda = \mathcal{O}(1\mu)$$

It follows that the second term within the curly brackets in Eq. (5) is several orders of magnitude smaller than the first term $\epsilon/\epsilon_0 \underline{E}$. It is therefore justifiable to neglect the second term and write Eq. (5) in the form

$$\nabla^2 \underline{E} + k_0^2 \frac{\epsilon}{\epsilon_0} \underline{E} = 0$$

The physical significance of replacing Eq. (5) by Eq. (10) is this. By discarding the term $\nabla \frac{1}{\epsilon} \nabla \epsilon \underline{E}$, we are neglecting any depolarization effects that may occur. This means that the wave retains the polarization it has at the source, which is evidenced by the fact that Eq. (10) can be reduced to a scalar equation by writing $\underline{E}(x)$ in the form

$$\underline{E}(\underline{x}) = \underline{e}_p \underline{u}(\underline{x}) \quad (11)$$

where \underline{e}_p is a unit vector in the direction of the initial polarization of the wave.⁴ Substituting Eq. (11) in Eq. (10), we find that $\underline{u}(\underline{x})$ satisfies the scalar wave equation,

$$\nabla^2 \underline{u} + k_0^2 \frac{\epsilon}{\epsilon_0} \underline{u} = 0 \quad (12)$$

This equation with the boundary condition on the initial surface, and the radiation condition at infinity, completely specifies $\underline{u}(\underline{x})$, from which we can then obtain the electromagnetic field vectors \underline{E} and \underline{H} .

If we write \underline{u} as the product of a factor $e^{ikn_0 z}$ that accounts for the rapid change in the phase of \underline{u} along the direction of propagation and a complex amplitude $A(\underline{x}, z)$, a further simplification of the problem results

$$[2ikn_0 \frac{\partial}{\partial z} + \nabla_T^2 + k^2(n^2(\underline{x}, z) - n_0^2)] A(\underline{x}, z) = - \frac{\partial^2 A(\underline{x}, z)}{\partial z^2} \quad (13)$$

where ∇_T^2 is the transverse Laplacian $\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$, and n_0 is a given constant which represents the refractive index of some uniform medium. At laser wavelengths the complex amplitude $A(\underline{x})$ varies much more rapidly transverse to the direction of propagation than it does along the direction of propagation. This enables us to make the paraxial approximation wherein the term on the right side of Eq. (13) is neglected (in the Russian literature this is called the parabolic approximation). So, the complex amplitude now satisfies

$$\left[i2kn_0 \frac{\partial}{\partial z} + \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + k^2(n^2(\underline{x}, z) - n_0^2) \right] A(\underline{x}, z) = 0 \quad (14)$$

For given initial data, i.e., values of the field at points on the initial surface, the propagation simulator must generate the corresponding field values at the terminal aperture such that Eq. (14) is satisfied. To do this we divide the medium into slabs defined by planes on which z is constant. In going from one slab to the next, we write $A(\underline{x}, z)$ in the form

$$A(\underline{x}, z) = e^{\Gamma(\underline{x}, z)} w(\underline{x}, z) \quad (15)$$

where $\Gamma(\underline{x}, z)$ is a phase function associated with the medium inhomogeneities

$$\Gamma(\underline{x}, z) = \frac{ik}{2} \int_{z_0}^z \left[n^2(\underline{x}, y, z') - n_o^2 \right] dz'. \quad (16)$$

The modified complex amplitude $w(\underline{x}, z)$ then satisfies the equation

$$\left[i2kn_o \frac{\partial}{\partial z} + e^{-\Gamma} v_T^2 e^{\Gamma} \right] w(\underline{x}, z) = 0 \quad (17)$$

with the initial condition

$$w(\underline{x}, y, 0) = u(\underline{x}, y, 0) \quad (18)$$

Physically, these equations approximate the propagation in the inhomogeneous medium by a two-step process at each z increment. First, we propagate the field $u(\underline{x})$ at $z - \Delta z/2$ to $z + \Delta z/2$, assuming that the intervening space is homogeneous. The effect of the inhomogeneities between $z - \Delta z/2$ and $z + \Delta z/2$ is then accounted for by multiplying this solution by the phase factor $\exp(\Gamma)$.

B. Adaptive Coordinates

To reduce the size of the mesh required to solve Eq. (17) numerically,

let us introduce an adaptive coordinate system defined by the transformation⁵

$$\zeta_1 = \frac{x/\rho_0}{N(z)} \quad (19)$$

$$\zeta_2 = \frac{y/\rho_0}{N(z)} \quad (20)$$

$$N(z) = \alpha^{-1/2} \left[\left(1 - \frac{z}{f} \right)^2 + \alpha^2 \left(\frac{z}{k\rho_0^2} \right)^2 \right]^{1/2} \quad (21)$$

$$\xi = \tan^{-1} \left[\frac{(1+\beta) \frac{z}{f}}{\beta^{1/2}} - 1 \right] \quad (22)$$

$$\beta^{1/2} = \alpha \frac{f}{k\rho_0^2} \quad (23)$$

where ρ_0 is a characteristic dimension of the beam at the initial surface (e.g., the e-folding radius of a gaussian beam), f is the distance to the focus, and α is a constant determined by the requirement that the solution be confined within the boundaries of the mesh at the focal plane. The choice $\alpha = 1$ yields a coordinate system that converges at a rate determined by the free-space diffraction of a gaussian beam having an e-folding radius ρ_0 .

When written in terms of the converging coordinate variables defined above, Eqs. (15) and (17) for the complex amplitude are replaced by the relations

$$w(x, y, z) = \hat{w}(\xi, \xi) \exp(\tilde{\Gamma}) v(\xi, \xi) \quad (24)$$

$$\hat{w}(\xi, \xi) = \left(\alpha^{1/2} N(z) \right)^{-1} \exp \left[\frac{i}{2} \left(\xi_1^2 + \xi_2^2 \right) \tan \xi \right] \quad (25)$$

$$\tilde{\Gamma} = \frac{ik}{2} \int_{z-\Delta z/2}^{z+\Delta z/2} dz' \left(n^2(x, y, z') - 1 \right) - \frac{i}{2} \left(\xi_1^2 + \xi_2^2 \right) \Delta \xi \quad (26)$$

$$\left[\frac{\partial}{\partial \xi} - \frac{i}{2} \exp(-\tilde{\Gamma}) \left(\frac{\partial^2}{\partial \xi_1^2} + \frac{\partial^2}{\partial \xi_2^2} \right) \exp(\tilde{\Gamma}) \right] v = 0 \quad (27)$$

where $\Delta\xi$ is the increment in ξ in going from $z-\Delta z/2$ to $z+\Delta z/2$. The initial condition for v is

$$v(\xi, \xi) = w(x, y, z)/\hat{w}(\xi, \xi) \quad (28)$$

To solve Eq. (27) we utilize the fact that for sufficiently small values of $\Delta\xi$ (i.e., Δz) the effect of the exponential factors $\exp(\pm \tilde{\Gamma})$ in this equation is small. Hence, we solve the simpler equation obtained when these factors are equated to unity

$$\left[\frac{\partial}{\partial\xi} - \frac{i}{2} \left(\frac{\partial^2}{\partial\xi_1^2} + \frac{\partial^2}{\partial\xi_2^2} \right) \right] v = 0 \quad (29)$$

We use a fast Fourier transform technique to solve Eq. (29). The basis of this approach is the fact that the solution of Eq. (29) can be expressed in the form of a discrete Fourier series

$$v(\xi, \xi) = \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} v_{mn}(\xi, t) \exp[i(p_m \xi_1 + q_n \xi_2)] \quad (30)$$

where the Fourier coefficients v_{mn} are determined from the initial data and Eq. (29) as follows. The initial values of v_{mn} are obtained by taking the discrete Fourier transform of the initial values of $v(\xi_i, \xi_i)$ over a mesh of points $\xi_1 = [\ell - (N/2)] \Delta\xi$, $\xi_2 = [j - (N/2)] \Delta\xi$ ($\ell, j = 0, 1, \dots, N-1$)

$$v_{mn}(\xi_i) = \frac{(-1)^{m+n}}{N^2} \sum_{\ell=0}^{N-1} \sum_{j=0}^{N-1} v\left(\left(\ell - \frac{N}{2}\right) \Delta\xi, \left(j - \frac{N}{2}\right) \Delta\xi, \xi_i\right) \exp\left[-\frac{i2\pi}{N} (m\ell + nj)\right] \quad (31)$$

The dependence of V_{mn} is then determined by substituting Eq. (30) in Eq. (31), which yields

$$\frac{\partial V_{mn}}{\partial \xi} + \frac{i}{2} \left(p_m^2 + q_n^2 \right) V_{mn} = 0 \quad (32)$$

from which it follows that

$$V_{mn}(\xi) = V_{mn}(\xi_i) \exp \left[- \frac{i(p_m^2 + q_n^2) \Delta \xi}{2} \right] \quad (33)$$

Finally, it can be shown that in order for the discrete Fourier series representation of v given in Eq. (30) to be real when v is real, the coefficients p_m and q_n must have the form

$$p_m = \frac{2\pi}{N \Delta \xi} \left(m - \frac{N}{2} \right) \quad (34)$$

$$q_n = \frac{2\pi}{N \Delta \xi} \left(n - \frac{N}{2} \right) \quad (35)$$

Hence, for discrete points $\xi_1 = (\ell - N/2) \Delta \xi$, $\xi_2 = (j - N/2) \Delta \xi$ ($\ell, j = 0, 1, \dots, N-1$)

$$v \left(\left(\ell - \frac{N}{2} \right) \Delta \xi, \left(j - \frac{N}{2} \right) \Delta \xi, \xi \right)$$

$$= (-1)^{\ell+j} \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} (-1)^{m+n} V_{mn}(\xi_i)$$

$$\exp \left[-i \hat{\beta} \left(\left(\frac{m - \frac{N}{2}}{N} \right)^2 + \left(\frac{n - \frac{N}{2}}{N} \right)^2 \right) + i \frac{2\pi}{N} (\ell m + j n) \right] \quad (36)$$

where $\hat{\beta} = 2\pi^2 \Delta\xi / (\Delta\xi)^2$. Note that v is simply $(-1)^{l+j}$ times the discrete Fourier transform of $(-1)^{m+n} v_{mn}(\xi)$.

The effect of the medium and the factor $\exp[(-i/2)(\xi_1^2 + \xi_2^2)\Delta\xi]$ introduced by the coordinate transformation is taken into account at each ξ step in the calculation by multiplying the value of v obtained in the previous step by the quantity $\exp(\tilde{\Gamma})$ defined in Eq. (26), i.e., the initial value inserted in Eq. (31) is $\exp(\tilde{\Gamma})$ times the value of v determined from the previous steps.

Using this adaptive coordinate algorithm we have been successful in our treatment of various realistic multimode fiber structures. Results are summarized in Section III.

C. Implementation of Lossy Outer Boundary.

It is believed that the field touching the outer boundary of the cladding region of a fiber structure will be attenuated due to radiation or absorption. To accomodate this situation in order to further improve our computer simulation, we have incorporated the presence of a lossy dielectric layer outside the cladding region in our computer program. An example of the index profile of a fiber is shown in Fig. 1:

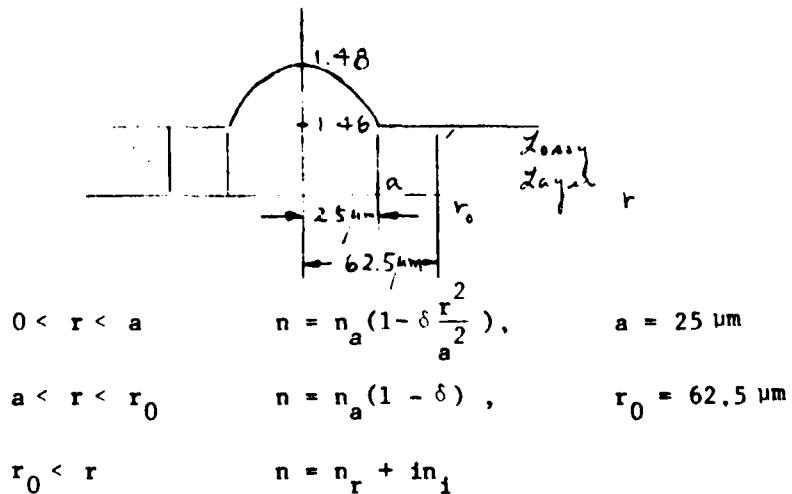


Figure 1: A Typical Index Profile with Lossy Outer Layer

Typical intensity patterns of beams propagating in a fiber with the index profile given by Fig. 1 are shown in Fig. 2. It was found that the propagation characteristics of the guided beams are not significantly affected by the presence of a lossy outer layer, except when the spot size of the beam is larger than the core diameter, as expected.

D. A Heuristic Approach in Obtaining the Reflection Coefficient.

One of the a'prior assumption in the development of the FFT scalar wave approach is that only paraxial rays are allowed and no reflection is permitted. This assumption enables us to develop an algorithm, thereby, we may obtain the propagating field by a forward stepping process as described earlier. As the field evolves from one z plane to the next $z + \Delta z$ plane, the averaged value of the refractive index as seen by the field may be different as illustrated in Fig. 3:

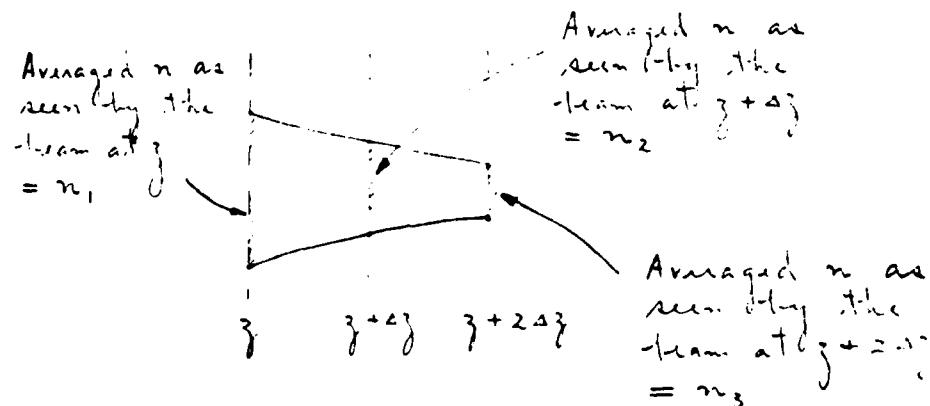
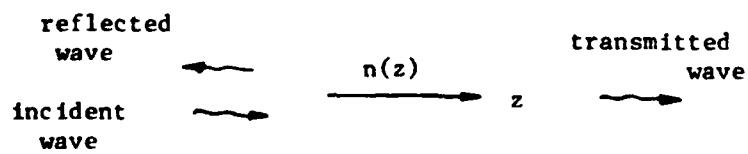


Figure 3: Illustration of the averaged n as seen by the beam.

In effect one may postulate that the wave is experiencing reflection in a medium with longitudinally slowly varying refractive index as shown below:



$n(z)$ is given by Fig. 3.

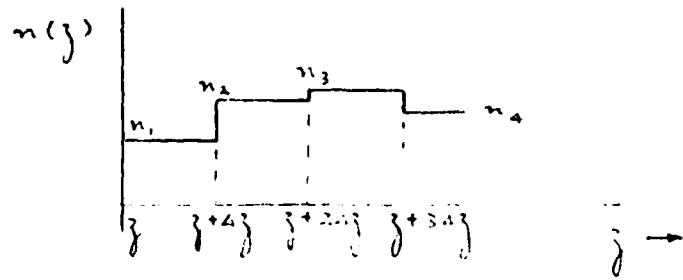


Figure 4: Equivalent Index Profile

The reflection coefficient for a plane wave propagating in this longitudinally non-uniform medium may be obtained according to a formula derived for the case of plane wave propagation in stratified layered medium:⁶

$$R(z) = - \exp [-is(z)] \int_z^\infty \gamma(z) \exp [is(z)] dz$$

$$s(z) = 2 \int_z^z \beta(z) dz \quad \beta(z) = k_0 n(z)$$

$$\gamma(z) = \frac{d\beta}{dz} / 2\beta$$

This is the heuristic approach that we shall use to calculate the reflection coefficient for waves in our multimode fiber structures.

III. Results

The algorithms detailed above have been implemented in our computer programs.

Results for the proposed tasks are given in the following:

(a) Effects of Step Index Gradient on the Propagation Characteristics.

The purpose of this study is to learn the effects of step index gradient on the propagation characteristics of waves in a multimode fiber guide. Let us introduce the following index profile:

$$n(r) = n_0 - \delta \left(\frac{r}{a}\right)^{2m} \quad (0 < r < a)$$

$$n(r) = n_0 - \delta \quad (a < r)$$

where n_0 , m and δ are given constants and a is the core radius of the fiber. For a typical parabolic index profile fiber, one has

$$n_0 = 1.48, \quad \delta = 0.02, \quad a = 25\mu\text{m}, \quad m = 1.$$

The constant δ must necessarily be small so that the depolarization effect may be ignored and the scalar wave approach may be justified. By varying m , the steepness of the index gradient may be varied as shown in Fig. 5. It should be kept in mind that even when the FFT technique is capable of handling steep index variations, the slope of the index profile must still be gentle enough so that the gradient term in the exact wave equation (Eq. (5)) may be ignored. We have carried out propagation calculation for the following specific cases: $n_0 = 1.48$, $\delta = 0.02$, $a = 25\mu\text{m}$, $m = 1, 4, 6, 10$. Higher m values mean steeper index gradient. As shown in Figs. 6, no computational difficulties were encountered for even the steepest case ($m = 10$) in which the index changes from $n = 1.48$ to $n = 1.46$ in $3\mu\text{m}$ distance for $50\mu\text{m}$ core diameter fiber. However, one should be aware that we are pushing the limit of validity for the scalar wave approach.

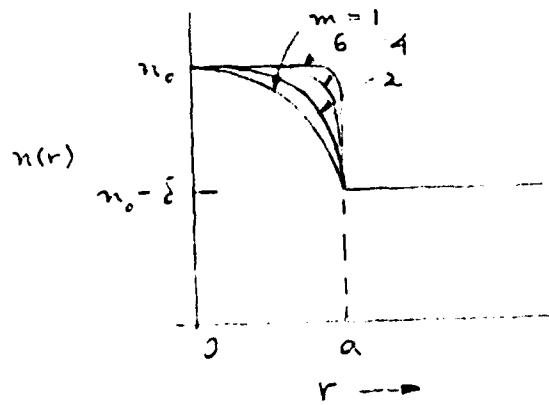


Fig. 5. Plot of $n(r)$ if $n(r) = n_0 - \delta (r/a)^{2m}$.

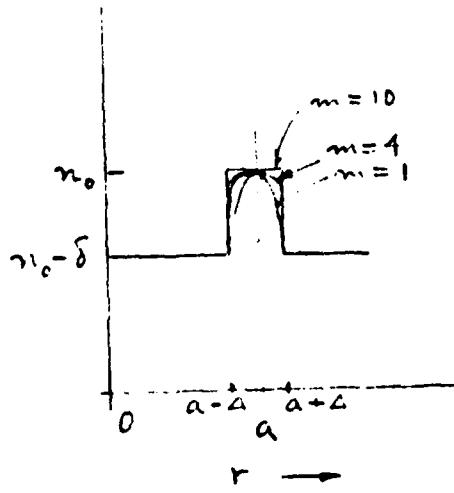


Fig. 7 Plot of $n = n_0 - \delta ((r-a)/\Delta)^{2m}$

From the results, it is of interest to note that as m increases from 1, i.e., as the index profile deviates from the parabolic profile, the beam profiles no longer remain to be of gaussian shapes, but take on ring-type structures. This implies that the phase front of the multimode beam is no longer a monotonic function of the radial distance but has become an oscillatory one.

We may conclude from these calculations that our program is capable of handling problems with steep index gradient. The index transition may occur in a distance as small as 4λ where λ is the free-space wavelength. The limiting factor apparently is the justification for the elimination of the depolarization term in Eq. (5).

(b) Beam Propagation in a Ring Fiber

A typical single-mode fiber has a core diameter of the order of $10\mu\text{m}$. Consequently it is more difficult to handle than a multi-mode fiber. An idea to enlarge the single-mode fiber has recently been put forth by Dr. L. Eyes of RADC. He suggested that perhaps a ring-type structure may support a single mode and yet possesses larger dimension than the usual solid-core fiber. This task was undertaken to investigate this possibility. Let us postulate that the index profile of a ring fiber takes the following

form
$$n = n_0 - \delta \left(\frac{r-a}{\Delta}\right)^{2m} \quad \text{for } a-\Delta < r < a+\Delta$$

$$n = n_0 - \delta \quad \text{for } r > a+\Delta$$

where n_0 , δ , a , Δ , and m are given constants. By increasing the m value, one may adjust the steepness of the index gradient as shown in Fig. 7. Two types of initial beam shapes will be studied: (1) a solid centered gaussian beam and (2) a hollow-centered donut beam. We wish to learn how well the ring fiber will confine these two types of beams. The field expression for a solid gaussian beam takes the form

$$u = e^{-\frac{\alpha}{2} \left(\frac{r}{a}\right)^2}$$

while the hollow-centered donut beam takes the form

$$u = e^{-\frac{\alpha}{2} \left[\left(\frac{r}{a}\right)^2 - 1 \right]}$$

where α and a are given constants. Results of our computation are shown in Figs. 8 and 9. By following the evolution of the beam intensities, one may determine how well the ring fiber is guiding the beam. It can be seen from these figures that the solid beams appear to be better confined than the hollow beams, although the spreading of the solid beam energy is quite noticeable. It also appears that simple insertion of beam energy in the high index region of the ring fiber does not insure good guidance of the beam energy. One may conclude from this preliminary study that neither solid Gaussian beams nor hollow-centered Gaussian beams correspond to the mode energy distribution of a single-mode in a ring fiber. One should first perform the classical modal analysis to obtain the mode pattern of the single mode and then use this mode pattern as the initial beam pattern for propagation down the ring fiber. It is believed that the use of ring-index fiber as large core single-mode fiber definitely possesses merit and should be studied further. What we have demonstrated with our present study is that our program is capable of handling this type of fibers.

(c) Fiber Couplers

One of the simplest type of light couplers is the fiber coupler. By placing two or more fibers in close proximity of each other light energy may transform from one to the other through the coupling effect. This coupling process is rather involved. The well-known coupled mode theory may be adequate for simple, single-mode structures such as slabs

with reasonable separations. But, when multi-mode complex structures such as the fiber couplers are involved, the coupled mode theory becomes grossly inadequate.* On the other hand, our FFT-scalar wave approach is uniquely qualified to deal with this fiber coupler problem. This is because this technique provides the evolution of beam field as it propagates down a complex multimode inhomogeneous fiber structure. Four types of fiber couplers have been studied:

Case 1 Coupling between two equal parabolic index fibers.

Two graded-index fibers are fused together longitudinally with separation d between their centers. The index profile for each fiber is given by

$$n(r_{1,2}) = n_0 - \delta \left(\frac{r_{1,2}}{a} \right)^2$$

where n_0 , δ , and a are given constants, and 1 or 2 refers, respectively, to #1 or #2 fiber. Typical values for a Corning or ITT graded index fibers are used:

$$n_0 = 1.48$$

$$\delta = 0.02$$

$$a = 25\mu\text{m}$$

*Recent advances by L. Eyes and P. Gianino of RADC using the extended boundary condition technique have shown that single mode couplers involving arbitrarily shaped uniform core guides can be successfully and accurately treated.

Various separation d were used. A gaussian beam represented by

$$u(x,y) = u_0 \exp \left\{ \left[- \left(x + \frac{d}{2} \right)^2 - y^2 \right] / w^2 \right\}$$

where $u(x,y)$ is the scalar wave function of the beam, and u_0 , w are given constants, is incident on one of the fibers. Results have been obtained for

$$w = 2.5\mu\text{m}, 5\mu\text{m}, 10\mu\text{m}$$

$$d = 8\mu\text{m}, 12\mu\text{m}, 16\mu\text{m}, 20\mu\text{m}.$$

The evolution of the beam along this structure is shown in Figs. 10-11. Displayed in Figs. 12 is the % power in one fiber as a function of longitudinal distance for various separations and initial beam sizes. For the situations considered above, many modes are excited. The coupling process is very involved as displayed in Fig. 10-11. It still appears that back and forth power exchange among the guides prevails. The complexity of the power exchange phenomenon for the multimode coupler re-emphasizes the importance of obtaining design data through analysis before the actual construction of fiber coupler.

Case 2 Coupling between two equal step-index fibers.

This coupler is identical to the previous one except step-index fibers were used. We shall approximate the index profile of a step index fiber by the following expression:

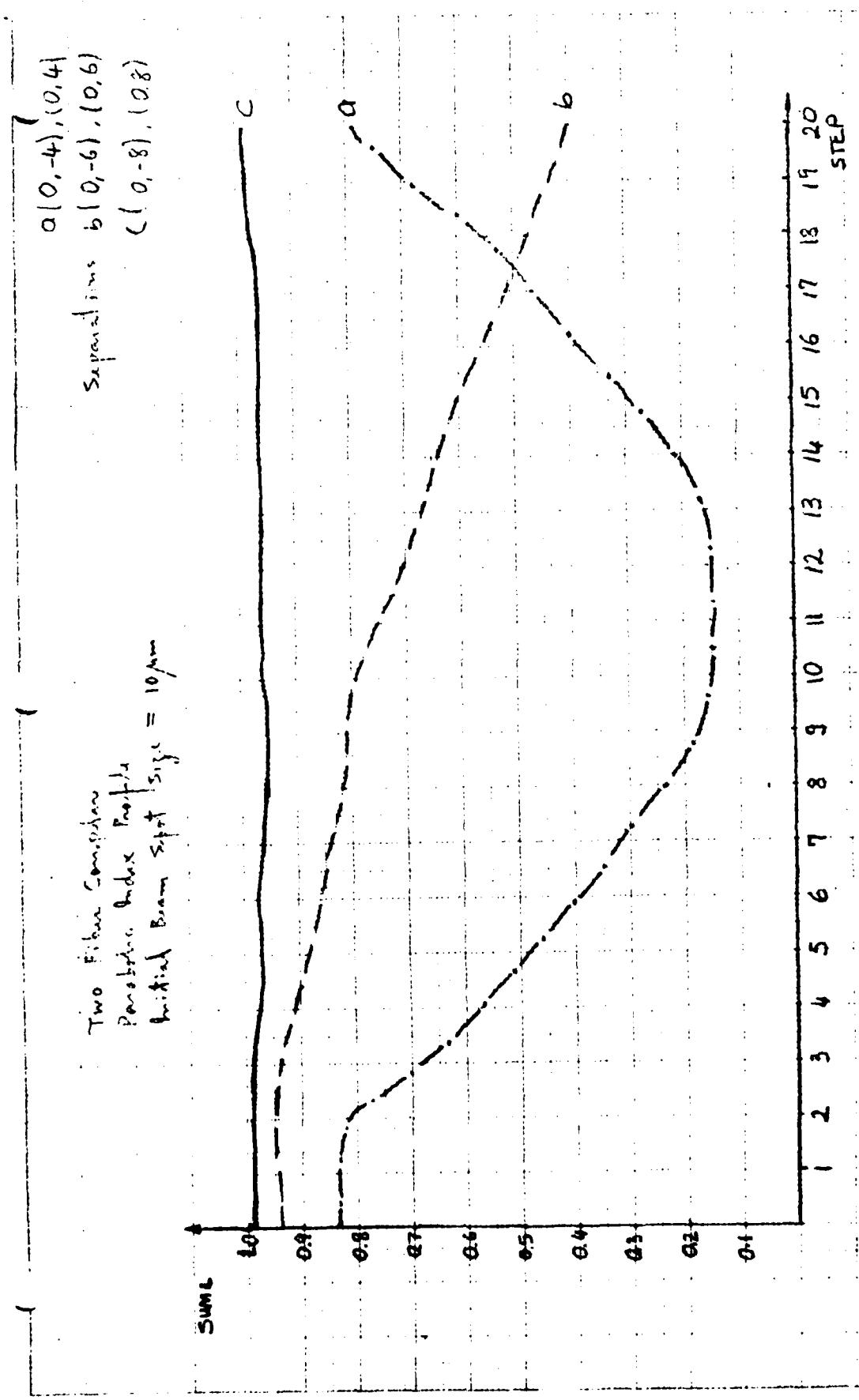


FIG 12 (a)

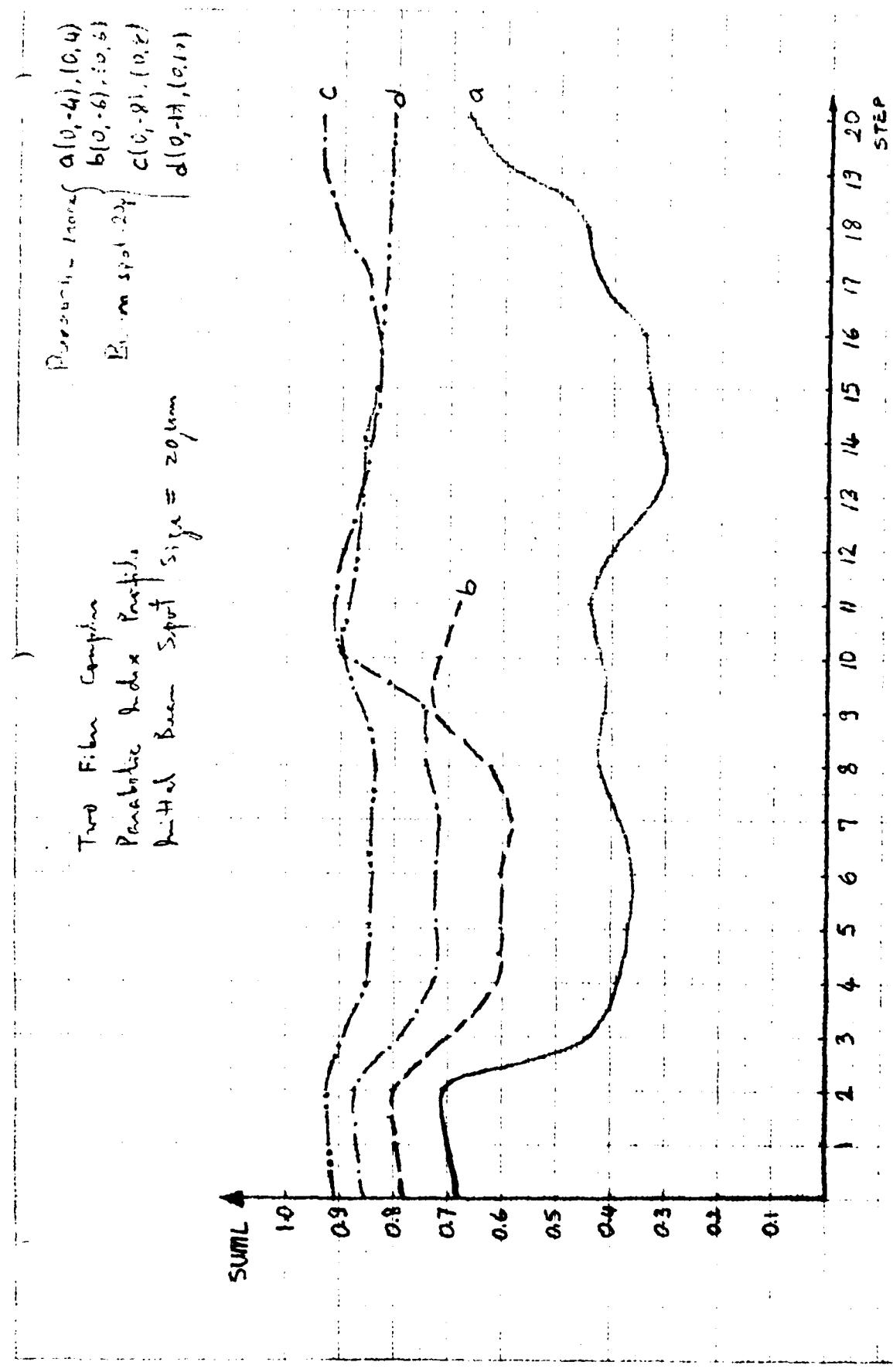


FIG 12 (b)

$$n(r_{1,2}) = n_0 - \delta \left(\frac{r_{1,2}}{a} \right)^{2m}$$

with $m = 4$. Again typical values for a Corning or ITT graded index fiber are used; i.e., $n_0 = 1.48$, $\delta = 0.02$, $a = 25\mu\text{m}$. Results have been obtained for

$$w = 2.5\mu\text{m}, 5\mu\text{m}$$

$$d = 8\mu\text{m}, 12\mu\text{m}, 16\mu\text{m}$$

where d is the separation distance and w is the beam waist radius.

Specific results are given in Figures 13. Displayed in Figs. 14 is the % power in one fiber as a function of longitudinal distance for various separations.

Case 3 Coupling between two unequal fibers.

It is of interest to learn, when two unequal size fibers are placed side by side, whether transfer of power would occur for the multimode case. The following fibers were used:

$$n(r_{1,2}) = n_0 - \delta \left(\frac{r_{1,2}}{a_{1,2}} \right)^{2m}$$

$$n_0 = 1.48 \quad m = 1, 4$$

$$\delta = 0.02$$

$$a_1 = 25\mu\text{m}$$

$$a_2 = 12.5\mu\text{m}$$

$$\text{Separation distance, } d = 8\mu\text{m}, 12\mu\text{m}$$

$$\text{Initial Beam Radius, } w = 5\mu\text{m}, 10\mu\text{m}$$

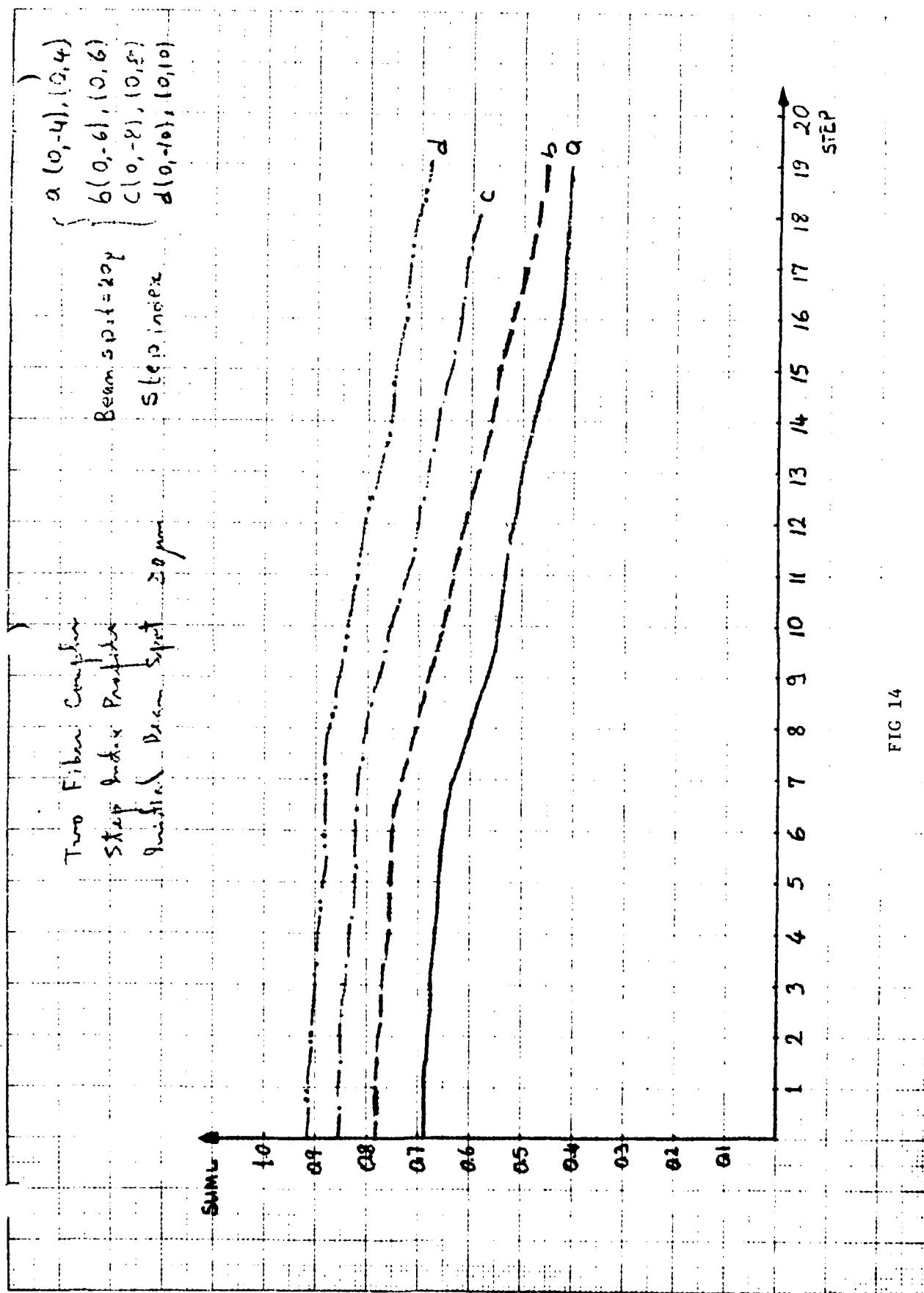


FIG 14

For the above chosen parameters, one may observe from Figs. 15-16 that only nominal coupling occurs between two unequal size fibers. In other words these structures are not efficient couplers.

Case 4 Coupling between more than two fibers.

As illustrations, we have considered two types of structures: Three equal size fibers located equal-distance from each other and three unequal size fibers arranged in a triangle shape. One fiber is initially illuminated, we wish to learn the power exchange characteristics in these two couplers. The following parameters were used:

$$n(r_{1,2,3}) = n_0 - \delta \left(\frac{r_{1,2,3}}{a_{1,2,3}} \right)^2$$

$$n_0 = 1.48$$

$$\delta = 0.02$$

$$a_1 = 25\mu\text{m}$$

$$a_2 = 25\mu\text{m}, 20\mu\text{m}$$

$$a_3 = 25\mu\text{m}, 15\mu\text{m}$$

$$\text{Separation distances, } d_1 = d_2 = d_3 = 12\mu\text{m}$$

$$\text{Initial beam radius, } w = 5\mu\text{m}$$

It can be seen from Fig.17 that when one of the fiber of 3 identical fibers, separated an equal distance from each other, is illuminated with a gaussian beam, power exchange takes place between two unilluminated fibers with the one illuminated fiber in a synchronous manner. In other words, each of the originally unilluminated fibers is receiving identical amount of power transfer from the illuminated one, as expected. On the other hand, the power exchange phenomenon for 3 non-identical fibers coupler is much more complicated as seen from Fig.18

Larger power transfer tends to take place among fibers with similar diameters.

Fiber coupler is one of the crucial components in any fiber optics system. As such, one must understand the detailed wave interaction phenomenon in this type of coupler so that correct design can be made. We have demonstrated the capability of our technique in dealing with this type of fiber structures. Systematic studies of different variations of fiber couplers may now be undertaken.

(d) Reflection Coefficient Calculations.

We have implemented the heuristic approach discussed earlier in the computer program to yield reflection coefficients for waves propagating in the various multimode fiber structures. Although it is difficult to justify the accuracy of the absolute values for the reflection coefficients obtained according to this algorithm, nevertheless, we feel that their relative values for different structures can be believed. This is because our heuristic approach took into consideration the fundamental characteristic of wave reflection: i.e., reflection occurs when discontinuity of the propagation medium or structure is experienced by the wave. The larger is the discontinuity, the larger will be the reflected energy.

Results for sample calculations for the reflection coefficients for various fiber structures are shown in Figs.19-20

IV. Conclusions and Recommendations

Support of this program has enabled the contractor to develop and perfect a computer program based on the scalar wave - FFT algorithm to study the propagation characteristics of guided waves in several important, practical fiber structures such as fibers with general index profiles (step index, parabolic index, ring index, etc), multi-channel fiber couplers, and fiber horns or tapers. These fiber structures may be made with commercially available fibers whose index variation may be as large as 1 - 2%.

We have implemented the adaptive coordinates and lossy outer mesh boundary schemes in our computer program. However, for most practical situations of interest in which the fiber core radius is about $25\mu\text{m}$, the cladding index is about 1 - 2% less than the core index, the spot size of the beam is less than $20\mu\text{m}$ and the free-space wavelength of the beam is larger than $0.6\mu\text{m}$, it is not necessary to implement the adaptive coordinates and lossy outer mesh boundary schemes. We also discovered that steep index gradient is not a hindrance for the program to produce accurate results as long as the scalar wave approach is still justifiable.

It is not unreasonable to ask the following question:

"Now that we have completed a beautiful program capable of producing propagation results for a variety of practical fiber structures, what can we do with it?"

The answer is "May be a lot!" Listed below are only a few of the problems that we can solve with this program:

- (1) Any single-mode or multimode weakly guiding fiber with arbitrary refractive index profile.

Our program provides the means to obtain the propagation characteristics of guided waves supported by this structure. The core of the structure may be circular, elliptical, rectangular, triangular or dumb-bell

shape with general index profile.⁸

- (2) Any fiber couplers composed of parallel strands of two or more of the above fibers. This is the only program which can provide the detailed coupling characteristics of this type of structure. Prior knowledge of coupling characteristics of a coupler is the key to successfully design and construct fiber couplers.
- (3) Any transition elements derived from the above fibers. Transition elements such as tapers, horns, or mode converters or branches can all be analyzed by our program.

Recommended Future Research

In addition to the important practical problems mentioned above that can be solved by our approach, it is worthwhile to look into the future and seek out problems of potential importance and interest. For example:

- (1) Nonlinear Fiber.

Very high intensity is achievable in fibers. One may wish to learn the propagation characteristics of waves in a fiber in which the induced nonlinearity of the medium plays a significant role. This problem may be solved by the scalar wave - FFT approach.

- (2) Mode Conversion in a Fiber Due to its statistically Varying Medium.

This problem may also be approached from the scalar wave - FFT point of view. It is known that the presence of frozen-in statistically varying medium contributes to the mode conversion phenomenon in a fiber.⁹ A systematic study of this problem will reveal the severity of this effect in changing the dispersion characteristics of beam in this fiber.

- (3) Large-size Single-Mode Fiber

The advantages of having large-size single-mode fiber are well-known. Ring-index fiber as proposed by L EYges of RADC appears

to be a promising one. Other type of fiber whose index variation may be radially unsymmetrical, such as, a layered fiber as shown in Fig. 21 may also be promising. Research should be encouraged on this type of fibers.

(4) Polarization Preserving Fibers

One of the main features of a weakly guiding fiber which can be analyzed by the scalar wave approach is that the guiding structure is polarization insensitive. However, for several important application areas, the polarization preserving characteristic of fiber is essential. Stress induced birefringent fiber or deformed core fiber may satisfy the needed requirement. However, to achieve the targeted isolation for the two orthogonal dominant modes, the required stress is excessively high for stress induced birefringent fiber and the loss is excessive for deformed-core fiber. We propose the use of layered dielectrics as shown in Fig. 21 to produce an equivalent birefringent effect to enable proper isolation for the two orthogonal dominant modes.¹⁰ Initial indication is very promising. It would therefore be very worthwhile to pursue this research.

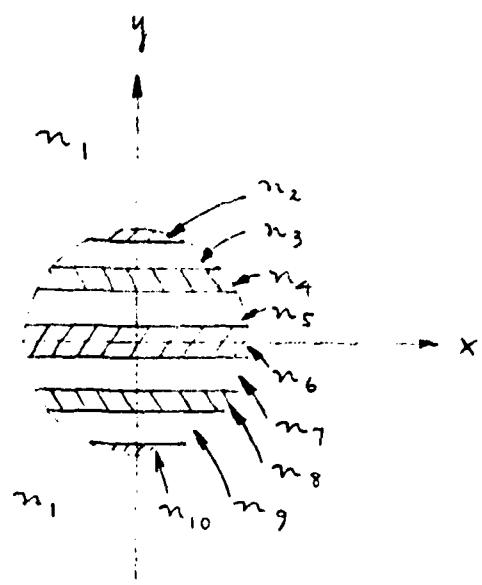


FIG. 21 Polarization Preserving Layered Fiber

Personnel:

The principal contributors of this contract have been:

C. Yeh	Senior Research Engineer
P. Barber	Research Engineer
F. Manshadi	Research Engineer

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9. M. Imai, T. Asakura and Y. Kinoshita, Opt. and Quan. Elec. 7, 95 (1975).
10. Patent Pending.

SAMPLE PROGRAM LISTINGS FOR THE CASE OF BEAM
PROPAGATION IN A STEP INDEX FIBER

CORE SIZE (DIAMETER) = 50 μ m
CORE INDEX = 1.48
CLADDING INDEX = 1.46
INITIAL GAUSSIAN BEAM SPOT SIZE (DIAMETER) = 40 μ m
WAVELENGTH = 0.8 μ m

(STEP INDEX CASE, m = 6)

LEVEL 2.3 - (JUNE 76)

OS/360 UTRAN H EXTENDED

DATE 00.19 1.57.04

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTOBLK(NONE)
SOURCE EBCDIC NOLIST NODECK OBJECT Q4A2 NOFORMAT GOSTMT NUKEY= ALC NOANSF NOTERM IOM FLA

```
C THIS PROGRAM CONTAINS: CORRECTED = 1834=4(JPTF1B,PRINTER,PEAK,SIZE,  
C HARY, GREYSC).  
C PROGRAM JPTF1B( INPUT, OUTPUT, =1-EJ,RA>E5=INPUT,TAPE6=OUTPUT,  
C + TAPE7=FILED),  
C  
C INPUT PARAMETERS  
C CARD 1 (15 FORMATS)  
C NCASES NUMBER OF CASES TO BE READ  
C CARD 2 (NAMELIST FORMAT-DEFAULT &  
C LAMBDA WAVE LENGTH (MICRONS)  
C RO 1/E POINT IN IRRADIANCE (MICRONS)  
C FR FIBER RADIUS (MICRONS)  
C NO REFRACTIVE INDEX  
C PCDRP PERCENT DROP AT R=FR OF 1/E  
C OUTRAD OUTER RADUS (MICRONS)  
C DX 42 SH SPACING (MICRONS)  
C NSTEPS NUMBER OF 2-STEP = 2MIN/NDZINC  
C NDZINC LENGTH OF 2-STEP = 2MIN/NDZINC  
C OUT DEVICE NUMBER FOR OUTPUT  
C GREY DEVICE NUMBER FOR GREY;  
C GREY IF TRUE INT IRRADIANCE PROFILE AT EACH STEP  
C PRAIST IF TRUE WRITE 2ND ADJENTS AT EACH STEP  
C PLTNS IF TRUE PLOT 2ND MOMENTS VS DISTANCE  
C ZLTMAX IF TRUE PLOT PEAK INTENSITY VS DISTANCE  
C PLTFLD IF TRUE PLOT FIELD IRRAD AT END OF PROPAGATION ALONG I  
C PLTFLW IF TRUE DO ABOVE PLOT AT EVERY STEP  
C GRID SIZE (32, 64, 128)  
C  
C SN 0002 COMMON /LCM2/KEFND(116384),SYN1231,CS(1128),ZTSQ(1128),PDSQ(1128)  
C * LCM2  
C LEVEL 2,REFND,SN,CS,ZTSQ,PUISQ,AM24RY,RADARY  
C DIMENSION WORK(25),W(13),INV(12),I(13),J(12),K(13)  
C DIMENSION X(13),Y(13),PCDRP(31),FR(31),REFCF(131)  
C COMMON /ARRAYS,V(32768)  
C COMMON /PARAM,ZINC,MESH,LAMDDA,HU,FR,NO,PCDRP,OUTRAD,DX,NSTEPS,  
C NDZINC,MESH,MSH502,P1,WAVEN,2XS1,VS,NS,ME,M,MSHPTS,  
C COMMON /PANPL/GREY,PA1ST,PL1ST,PLIMAX,LAST,IGREY  
C *PLTFL,PLTFL  
C  
C SN 0003 REAL LAMBDA,NO,NO50,N2  
C LOGICAL PGREY,PA1ST,PLTNS,PLTMAX,PLTFLD,PLTFL,PLTFL,PLTFL  
C SN 0004  
C SN 0005  
C SN 0006  
C SN 0007  
C  
C SN 0008  
C SN 0009  
C SN 0010  
C SN 0011  
C  
C DATA ICNTCS/1/.ICNT/0/  
C  
C NAMELIST /DEFAULT/LAMBDA,RO,NO,PCDRP,DDJTRAD,DX,NSTEPS,NDZINC,  
C * IOUT,IGREY,PGREY,PA1ST,PL1ST,PLIMAX,PLTFLD,PLTFL,PLTFL  
C
```

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LEVEL 2--J (JUNE 78)      MAIN      OS/360  COBOL H EXTENDED      DATE 80-19  21-57-04
ISN 0012      READ(5,1000) NCASES      OPTF18
ISN 0013      WRITE(6,1000) NCASES      OPTF18
ISN 0014      READ(5,11) XBX,YB      TEMP
ISN 0015      WRITE(6,11) XBX,YB      TEMP
ISN 0016      READ(5,1000) NF1A      TEMP
ISN 0017      WRITE(6,1000) NF1A      TEMP
ISN 0018      READ(5,12) X0(1,1),Y0(1,1),I=1,N=13)      TEMP
ISN 0019      WRITE(6,12) X0(1,1),Y0(1,1),I=1,N=13)      TEMP
ISN 0020      READ(5,13)(PCDRPA(K),K=1,NF1A)      TEMP
ISN 0021      READ(5,11)(PCDRPA(K),K=1,NF1B)      TEMP
ISN 0022      WRITE(6,13)(PCDRPA(K),K=1,NF1B)      TEMP
ISN 0023      WRITE(6,11)(PCDRPA(K),K=1,NF1B)      TEMP
ISN 0024      1000      TEMP
ISN 0025      FORMAT(1X,12)      TEMP
ISN 0026      11      FORMAT(1X,F4.1,1X,F4.1)      TEMP
ISN 0027      12      FORMAT(1X,F12.5,1X,F12.5)      TEMP
ISN 0028      C1      READ(5,FAULT)      OPTF18
ISN 0029      C      MM(1)= 7      OPTF18
ISN 0030      MM(2)= 7      OPTF18
ISN 0031      MM(3)= 0      OPTF18
ISN 0032      ICF=0      OPTF18
ISN 0033      FLAG=NO/ABS(IND)      OPTF18
ISN 0034      NO=ABS(IND)      OPTF18
ISN 0035      IF FLAG='LT' .0.  WRITE(6,OUT*2050)      OPTF18
ISN 0036      2050      FORMAT('0.47M THE REFRACTIVE (N2) IS A CONSTANT EQUAL TO N3')      OPTF18
ISN 0037      PCDRP=PCDRPA(1)      TEMP
ISN 0038      PCDRPA(1)      TEMP
ISN 0039      2050      FDRPA(1)      TEMP
ISN 0040      WRITE(10,1,DEFAULT)      TEMP
ISN 0041      C      CALCULATE CONSTANTS      TEMP
ISN 0042      NESH2=2*MESH      TEMP
ISN 0043      NESH2SQ=2*MESHSQ      TEMP
ISN 0044      RN2NO=0.02*PCDRP*FA**2      TEMP
ISN 0045      ZMIN=P1/(L2*SORTRN2NO))      TEMP
ISN 0046      DZINC=ZMIN/NDZINC      TEMP
ISN 0047      DXS1=DZINC/ZMIN      TEMP
ISN 0048      DXS1H=DXS1/2.      TEMP
ISN 0049      DZE=D/R0      TEMP
ISN 0050      WAVENM=2*pi/LAMBDA      TEMP
ISN 0051      BETHT=(2**2*MINDXS1/(WAVENM))*(PI/(WESHDZET*RO))**2      TEMP
ISN 0052      FTNST=(1.-1./MESH)*PI      TEMP
ISN 0053      XYME SH/2.      TEMP
ISN 0054      RDRM=(UJTHAD/RO)**2      TEMP
ISN 0055      N2=ND*RN2NO      TEMP
ISN 0056      REFCC=N2*RU**2/2      TEMP
ISN 0057      ALPHA=2.*ZMIN/(PI*WAVENM*NJ*4*U**2)      TEMP
ISN 0058      NDSO=NJ*2      TEMP
ISN 0059      K SINDDXS1      TEMP
ISN 0060      LAS=.FALSE.      TEMP
ISN 0061      IF(FLAG.LT.0.)  REFCF=0.      TEMP

```

```

LEVEL 2.3.0 (JUNE 76)      MAIN      DS/350  - LTRAN II EXTENDED      DATE 80.199/. 57.04
1SN 0063      CALLPH=GREY,DR,PHASIT,DR,P-IMST,2-TMAX,DR,PLTFLD,DR,PLTLE      OPTF1B
C
C      WRITE THE IMPORTANT CALCULATED PARAMETERS      OPTF1B
C
C      0064      2000 FORTRAN 2000) ZMIN,DZINC,RN2NU,ALPHA      OPTF1B
1SN 0065      2000 FORTRAN 2000) ZMIN,DZINC,RN2NU,714(C4NS,72H DZINC,710.4.1X,      OPTF1B
C      7MICRONS,79H RN2ND,710.3.1X,1.7MICRONS,710.5.1X,      OPTF1B
C      QH ALPHA,710.5.1X,      OPTF1B
C
C      CALCULATE NECESSARY ARRAYS      OPTF1B
C
C      MP=6      OPTF1B
1SN 0066      DO 800 K=1,MP      OPTF1B
1SN 0067      X0(K)=X0(K)+XY0      OPTF1B
1SN 0068      Y0(K)=Y0(K)+XY0      OPTF1B
1SN 0069      MX=2*MP      OPTF1B
1SN 0070      MP=6      OPTF1B
1SN 0071      REFCA(1)=Y0*.02*PCORPA(K)*(Z0/2+X0(K))*MX/2.      TEMP
1SN 0072      600      TEMP
1SN 0073      DO 100 K=1,MESH      TEMP
1SN 0074      RK=K-1      TEMP
1SN 0075      ARG=FCNST*RK      TEMP
1SN 0076      CS(K)=COS(ARG)      TEMP
1SN 0077      SN(K)=SIN(ARG)      TEMP
1SN 0078      ZTSQ(1)=((RK-XY0)*DZET)*02      TEMP
1SN 0079      PQSQ(1)=((RK-XY0+5)*02      TEMP
1SN 0080      100      TEMP
C
C      SET UP REFRACTIVE INDEX ARRAY      TEMP
C
C      M=0      OPTF1B
1SN 0081      DO 120 J=1,MESH      OPTF1B
1SN 0082      DC 120 I=1,MESH      OPTF1B
1SN 0083      N=M+1      TEMP
1SN 0084      TMPNDX=0.      TEMP
1SN 0085      DO 110 K=1,MP      TEMP
1SN 0086      Z1=((J-1-Y0(K))*DZET)*02      TEMP
1SN 0087      Z2=((I-1-Y0(K))*DZET)*02      TEMP
1SN 0088      RADZ1=Z2      TEMP
1SN 0089      TMPNDX=MAX((TMPNDX,(NO-REFCA(K)*(RADZ1*MP)))      TEMP
1SN 0090      CREF=NUSL*0.01*PCORPA(K)      TEMP
1SN 0091      REFNDX=MAX((TMPNDX,CREF))      TEMP
1SN 0092      CONTINUE      TEMP
1SN 0093      110      TEMP
C      IF((I-EQ.52.04-.1.EQ.58) WRITE(11,11) REFNDX(M)      TEMP
C      120      TEMP
C      CONTINUE      TEMP
1SN 0094      111      TEMP
1SN 0095      IF(MESHNE.128) GO TO 10      TEMP
1SN 0096      M=MESH/4      TEMP
1SN 0097      NF=M+MESH/2-1      TEMP
1SN 0098      NS=M      TEMP
1SN 0099      NF=NF      TEMP
1SN 0100      DO 120 GO TO 40      TEMP
1SN 0101      10      TEMP
1SN 0102      M=MESH/4      TEMP
1SN 0103      NS=1      TEMP
1SN 0104      NS=1      TEMP
1SN 0105      NF=MESH      TEMP

```


LEVEL 2.3. (JUNE 78) MAIN 05/30 JTRAN H EXTENDED DATE 80/199 1.57.04

```

ISN 0150      V(K)=VREAR-VIAI
ISN 0151      V(KP1)=VIAI+VRAI
ISN 0152      CONTINUE
      C      DO TRANS=JRM
      C      CALL HARM(V,MM,INV,S,1,IFERR)
      C      SOLVE FIRST UNDER ODE
      C      K=-1
      C      DO 150 J=1,MESH
      C      PHI1=3*TAH*POSQ(J)
      C      DO 150 I=1,MESH
      C      K=K+2
      C      PHI1=TAH*POSQ(I)
      C      PHI2=2*TAH*POSQ(I)
      C      V=V(KP1)
      C      Y=V(KP1)
      C      ANG=(PHI1+PHI2)
      C      SANG=COS(ANG)
      C      SANG=SIN(ANG)
      C      V(K)=(V*ANG-VI*SANG)
      C      V(KP1)=VY*VANG+VSECANG)
      C      150 CONTINUE
      C      DO INVERSE TRANSFORM
      C      CALL HARM(V,MM,INV,S,-1,IFERR)
      C      RECONDITION V BECAUSE OF TRANSFORM
      C      K=-1
      C      DO 200 J=1,MESH
      C      SN=SN(J)
      C      CS=CS(J)
      C      R21=(J-KYD-1)*DZET)**2
      C      CLOSS=2,
      C      DO 200 I=1,MESH
      C      R72=(I-KYD-1)*DZET)**2
      C      RAD=B21*DZET
      C      CM=1
      C      FF(3340,GF,HAJNRM) CM=EXP(-C-USS*(RAD-RADNRM))
      C      K=K+2
      C      PHI2=K
      C      SN1=SN(1)
      C      CS1=CS(1)
      C      A=CSJ*CS1-SNJ*SNI
      C      A=CSJ*SN1+SNJ*CSI
      C      V=V(K)*CM
      C      V1=V(KP1)*CM
      C      V(K)=VREAR-VIAI
      C      V(KP1)=VRAI+VIAI
      C      CONTINUE
      200

```


LEVEL 2.3. (JUNE 79) OS/360 JETRAN H EXTENDED DATE 60.195 1.57.07

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(26) SIZE(MAX) AUTOBL(NONE) SOURCE EBCDIC NOLIST NODECK OBJECT NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IBM FLA

```
ISN 0002          SUBROUTINE PEAK(MAX)
ISN 0003          COMMON /LCM2/REFNDX(16384),N(123),CS(129),ZTS0(128),PQ50(128)
ISN 0004          *      AMPARY(16384),RADARY(16384)
ISN 0005          LEVEL 2,REFNDX,SNCS,ZTSQ,PJSU,A424M,RADARY
ISN 0006          COMMON /ARRAYS/V(32768)
ISN 0007          NOZINC,WEHSHQ,MESH502,PI,WAVEN4,2421,NS,WF,WF,MSH515
ISN 0008          L=0
ISN 0009          SUM=0
ISN 0010          SUMR=0
ISN 0011          VMAX=0
ISN 0012          DO 10 K=1,MSHSQ2,2
ISN 0013          VR(V,K)
ISN 0014          L=L+1
ISN 0015          KP=K+1
ISN 0016          VI=V(K1)
ISN 0017          VRAD=V16*2+VR*2
ISN 0018          IF(K,LE,WEHSHQ) SUML=SUML+VRAD
ISN 0019          IF(K,GE,WEHSHQ) SUMR=SUMR+VRAD
ISN 0020          VMAX=AMAX(VMAX,VRAD)
ISN 0021          RADAR(VL,VRAD)
ISN 0022          CONTINUE
ISN 0023          TOT=SUML+SUMR
ISN 0024          SUML=SUML/TOT
ISN 0025          SUMR=SUM3/TOT
ISN 0026          WRITE(6,2000) SUML,SUMR
ISN 0027          FORMAT(1,IX,7HSJVR = ,E14.7,JX,7HSJVR = ,E14.7,I)
ISN 0028          RETURN
ISN 0029          END
```

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(26) SIZE(MAX) AUTOBL(NONE)

OPTIONS IN EFFECT: SOURCE EBCDIC NOLIST NODECK OBJECT NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IBM FLA

STATISTICS: SOURCE STATEMENTS = 26, PROGRAM SIZE = 540, SUBPROGRAM NAME = PEAK

STATISTICS: NO DIAGNOSTICS GENERATED

***** END OF COMPIRATION *****

112K BYTES OF CORE NOT USED

REQUESTED OPTIONS:

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INTELLIGENT SYSTEMS IN ELECTRONICS

```

SOURCE EBCDIC NULIST NODECK OBJECT NO4AP NOFORMAT GOSTM NOXHEF ALL NUANSF NOTFRA 164 FLA
SUBROUTINE SIZE(LARRAY,MESH,MESH,X0,Y0,X2,Y2) SIZE
  DIMENSION ARRAY(MESH,MESH)
  MID=MESH/2+1
  SUM1=0.
  SUM2=0.
  SUM3=0.
  SUM4=0.
  SUM5=0.
  DO 10 N=1,MESH
  RN=N*MC
  RN5=RN*RN*2
  DO 10 M=1,MESH
  RM=M*MD
  RMS=RN*RM*2
  ARMN=ARAY(M,N)
  SUM1=SUM1+RM*ARMN
  SUM2=SUM2+RN*ARMN
  SUM3=SUM3+RN*ARMN
  SUM4=SUM4+RN*ARMN
  SUM5=SUM5+ARMN
  10  CONTINUE
  SNORM=1./SUM5
  X0=SNORM*SUM2
  Y0=SNORM*SUM1
  X2=SNORM*SUM4-X0*42
  Y2=SNORM*SUM3-Y0*42
  X2=SQRT(X2)
  Y2=SQRT(Y2)
  RETURN
  END

OPTIONS IN EFFECT NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBL(NONE)
OPTIONS IN EFFECT SOURCE EBCDIC NULIST NODECK OBJECT NO4AP NOFORMAT GOSTM NOXHEF ALC NOANSF NOTFRA 164 FLA
STATISTICS* SOURCE STATEMENTS = 30. PROGRAM SIZE = 786. SUBPROGRAM NAME = SIZE
STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILEATION *****

```

LEVEL 203-J (JUNE 78) OS/350 UNTRAN H EXTENDED DATE 00-19 1-57-09

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME (MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTOBL (NONE)
OPTIONS IN EFFECT: SOURCE EBCDIC NOLIST NODECK OBJECT NO44P NOFORMAT COSTNT NOXREF ALC NOANSF NOTFR4 10M FL/
INITIAL ARRAYS
ISN 0002 /ARRAYS/V(32768)
ISN 0003 COMMON /PARAHDZINC, MESH, LA4BDA, 4J, F4, NO, PCDRP, OUTRAD, DX, VSTEP, PARAM
ISN 0004 * NOZINC, MESHQ, MSHSQ, NS, MF, NS, MSHPS, PARAM
ISN 0005 COMMON /PARNTL, PGREY, PMAST, PLTST, PLTMX, LAST, IOUT, IGREY
* PLTFLO, PLTFILE
* LOGICAL PGREY, PMAST, PLTST, PLTMX, LAST
LOGICAL PLTFLO, PLTFILE
REAL LAMBDA, NO
DATA PGREY, PMAST, PLTMAX, LAST, IOUT, IGREY /20, TRUE, 20, FALSE/
ISN 0006 DATA PGREY 20/
* FALSE, 20/
ISN 0007 DATA LAMBDA, RD, NO, PCDRP, OUTRAD, DX, VSTEP, NOZINC
* /J, 80, 20, 0, 1.48, 1.4E+0, 0.01, 0.01, 0.5/
DATA PLTFLO, PLTFILE, PMAST, PLTST, PLTMX, LAST, IOUT, IGREY /
DATA MESH/128, PI/314159265/
END
ISN 0012
ISN 0013

OPTIONS IN EFFECT: NAME (MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTOBL (NONE)

OPTIONS IN EFFECT: SOURCE EBCDIC NOLIST NODECK OBJECT NO44P NOFORMAT COSTNT NOXREF ALC NOANSF NOTFR4 10M FL
OPTIONS IN EFFECT: SOURCE EBCDIC NOLIST NODECK OBJECT NO44P NOFORMAT COSTNT NOXREF ALC NOANSF NOTFR4 10M FL
* STATISTICS* SOURCE STATEMENTS = 12. PROGRAM SIZE = 0. SUBPROGRAM NAME = ARRAYS
* STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
112K BYTES OF CORE NOT USED

LEVEL 2.3.0 (JUNE 76)

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(12) LINECOUNT(56) SIZE(MAX) AUTOOR(NONE) SOURCE EBCDIC NOLIST NODECK OBJECT 40442 NDFORMAT GOSTAT NORME ALC NDANSF NOTERM IBM =L/

SUBROUTINE HARM

PURPOSE DISCRETE COMP-EX FOURIER TRANSFORMS ON A COMPLEX THREE DIMENSIONAL ARRAY

USAGE

CALL HARM (A,M,INV,S,IFSET,IFERR)

DESCRIPTION OF PARAMETERS

A - AS INPUT. A COMPLEX, 3-DIMENSIONAL ARRAY TO BE TRANSFORMED. THE REAL PART OF ALL (1,2,13) IS STORED IN VECTOR FASHION IN A CELL WITH INDEX $2*(13*N1 + 12*N2 + 11) + 1$ WHERE $N1 = 2*N1(1)$, $N2 = 2*N2(1)$ AND $N1 = 0, 1, \dots, N1-1$ ETC. THE IMAGINARY PART IS IN THE CELL IMMEDIATELY FOLLOWING. NOTE THAT THE SUBSCRIPT i INCREASES MOST RAPIDLY AND $i = 1$ INCREASES LEAST RAPIDLY. AS INPUT A COMPLEX FOURIER TRANSFORM. THE NUMBER OF CORE LOCATIONS OF ARRAY A IS $2*(N1*2*N2*N3)$.

M - A THREE CELL VECTOR WHICH DETERMINES THE SIZES OF THE 3 DIMENSIONS OF THE ARRAY A. THE SIZE $N1$ OF THE 1 DIMENSION OF A IS $2^M(1)$, $i = 1, 2, 3$.

INV - A VECTOR WORK AREA FOR SIT AND INDEX MANIPULATION. IF DIMENSION ONE = 2*J+1 OF THE QUANTITY $M(N1,N2,N3)$.

S - A VECTOR WORK AREA FOR SINE TABLES WITH DIMENSION THE SAME AS INV.

IFSET - AN OPTION PARAMETER WITH THE FOLLOWING SETTINGS
0 - SET U2 SINE AND INV TABLES ONLY
1 - SET U2 SINE AND INV TABLES ONLY AND CALCULATE FOURIER TRANSFORM
-1 - SET UP SINE AND INV TABLES ONLY AND CALCULATE INVERSE FOURIER TRANSFORM (FOR THE MEANING OF INVERSE SEE THE EQUATIONS UNDER SET 1,2,3 BELOW)
2 - CALCULATE FOURIER TRANSFORM ONLY (ASSUME SINE AND INV TABLES EXIST)

-2 - ERROR INDICATOR. WHEN IFSET IS 0, 1, 2
(ASSUME SINE AND INV TABLES EXIST)
CALCULATE INVERSE FOURIER TRANSFORM ONLY
-2 - ERROR INDICATOR. WHEN IFSET IS 0, 1, 2
IFERR 2 = 1 MEANS THE MAXIMUM WH. IS GREATER THAN $20 \cdot 1 \cdot 2 \cdot 3$ WHEN IFSET IS 2, -2 MEANS THAT THE SINE AND INV TABLES ARE NOT LARGE ENOUGH OR HAVE NOT BEEN COMPUTED
IF ON RETURN IFERR = 3 THEN NONE OF THE ABOVE CONDITIONS ARE PRESENT

OS/360 JETIAN M EXTENDED

DATE 00.199 1.57.09

LEVEL 2.3.0 (JUNE 76)

05/360 J474N H EXTENDED

DATE 00.199 .057.09

REMARKS
 THIS SUBROUTINE IS TO BE USED FOR COMPLEX, 3-DIMENSIONAL
 ARRAYS IN WHICH EACH DIMENSION IS A POWER OF 2. THE
 MAXIMUM N(1) MUST NOT BE LESS THAN 3 OR GREATER THAN 20.
 I = 1,2,3

SUBROUTINES AND FUNCTIONS REQUIRED
 NONE

METHOD
 FOR IFSET = +1, OR +2, THE FOURIER TRANSFORM OF COMPLEX
 ARRAY A IS OBTAINED.

$$x(j_1, j_2, j_3) = \sum_{k_1=0}^{N_1-1} \sum_{k_2=0}^{N_2-1} \sum_{k_3=0}^{N_3-1} A(k_1, k_2, k_3) \cdot w_1^{j_1 k_1} \cdot w_2^{j_2 k_2} \cdot w_3^{j_3 k_3}$$

WHERE w_i IS THE N(i) ROOT OF UNITY AND $L_i = k_i + j_i$.
 $L_2 = k_2 + j_2, L_3 = k_3 + j_3$

FOR IFSET = -1, OR -2, THE INVERSE FOURIER TRANSFORM A OF
 COMPLEX ARRAY X IS OBTAINED.

$$A(k_1, k_2, k_3) = \sum_{j_1=0}^{N_1-1} \sum_{j_2=0}^{N_2-1} \sum_{j_3=0}^{N_3-1} x(j_1, j_2, j_3) \cdot w_1^{-L_1 j_1} \cdot w_2^{-L_2 j_2} \cdot w_3^{-L_3 j_3}$$

SEE J. B. COOLEY AND J. W. TUKEY, "AN ALGORITHM FOR THE
 MACHINE CALCULATION OF COMPLEX FOURIER SERIES",
 MATHEMATICS OF COMPUTATION, VOL. 19 (APR. 1965), P. 297.

SUBROUTINE HARM(M, N, INVS, IFSET, IFE2A),
 DIMENSION A(1), INV(1), S(1), V(3), V(3), NP(3), W(2), W(2), W(2), W(2)
 EQUIVALENCE (IN, N(1)), (IN2, N(2)), (IN3, N(3))
 10 IF(ABS(IFSET) - 1) 900, 903, 12
 12 M1=MAX0(M(1), M(2), M(3)) - 2
 ROOT2 = SQRT(2.)
 14 (M1-M1) 14, 14, 13
 13 IF(IFSET = 1)
 15 RETURN
 16 IF(IFSET = 0)
 17 M1=M(1)
 18 M2=M(2)
 19 M3=M(3)
 20 N1=2*NP(1)
 21 N2=2*NP(2)
 22 N3=2*NP(3)
 23 IF(IFSET = 1) 16, 16, 20
 24 N1=NP(1)
 25 N2=NP(2)
 26 N3=NP(3)

LEVEL 2. J (JUNE 78) HARM OS/360 - J2TRAN H EXTENDED DATE 80.19 . 21.57.09

```

ISN 0020 FN = NX HARM HARM1070
ISN 0021 DO 19 L = 1,NX HARM1080
ISN 0022 A(2*1-1) = A(2*1-1)/FN HARM1090
ISN 0023 19 A(2*1) = -A(2*1)/FN HARM1100
ISN 0024 20 NP(1,1)=1*2 HARM1120
ISN 0025 NP(2,1)=NP(1,1)*N2 HARM1130
ISN 0026 NP(3,1)=NP(2,1)*N3 HARM1140
ISN 0027 DO 250 L=1,3 HARM1150
ISN 0028 LL = NP(3,1)-NP(1,0) HARM1160
ISN 0029 LL1 = LL+1 HARM1170
ISN 0030 M = M(1,0) HARM1180
ISN 0031 IF (411250,250,30 HARM1190
ISN 0032 30 IDIF=NP(1,0) HARM1200
ISN 0033 K11=NP(1,0) HARM1210
ISN 0034 NEV = 2*(M1/2) HARM1220
ISN 0035 IF (41 - 4EV )60,60,40 HARM1230
C M IS 000, DO L=1 CASE HARM1240
C 40 K11=NP(1,2) HARM1250
ISN 0036 KLL=KBL1-2 HARM1260
ISN 0037 DO 50 I=1, IL1,1IDIF HARM1270
ISN 0038 KLAST=KL+1 HARM1280
ISN 0039 DO 50 K=1,KLAST,2 HARM1290
ISN 0040 KDL=K+KBL1 HARM1300
ISN 0041 C DO ONE STEP WITH L=1, J=0 HARM1310
C C C A(K)=A(K)+A(KD) HARM1320
C C C A(KD)=A(K)-A(KD) HARM1330
ISN 0042 T=A(KD) HARM1340
ISN 0043 A(KD)=A(K)-T HARM1350
ISN 0044 A(K)=A(K)+T HARM1360
ISN 0045 T=A(KD) HARM1370
ISN 0046 A(KD)=A(K+1)-T HARM1380
ISN 0047 50 A(K+1)=A(K+1)+T HARM1390
ISN 0048 IF (41 - 1)250,52 HARM1400
ISN 0049 52 LFIRST =3 HARM1410
C DEF - JLAST = 2*(L-2) -1 HARM1420
ISN 0050 JLAST=1 HARM1430
ISN 0051 GOT3 70 HARM1440
C M IS EVEN HARM1450
ISN 0052 60 LFIRST = 2 HARM1460
ISN 0053 JLAST=0 HARM1490
ISN 0054 70 DO 240 L=L,FIRST,MI+2 HARM1500
ISN 0055 JDOIF=BIT HARM1520
ISN 0056 K11=KBL1/4 HARM1530
ISN 0057 KL=K311-2 HARM1540
C DO F3W J=0 HARM1550
ISN 0058 DO 80 I=1,IL1,1IDIF HARM1560
ISN 0059 KLAST=1+KL HARM1570
ISN 0060 DO 80 K=1,KLAST,2 HARM1580
HARM1600

```

LEVEL 2. 1 (JUNE 78) HARM OS/360 - FORTRAN H EXTENDED DATE 80-15 21-57-09

```

      K=K+KBIT
      K2=K1+KBIT
      K3=K2+KBIT
      C
      DO TWO STEPS WITH J=0
      A(K1)=A(K1)+A(K2)
      A(K2)=A(K1)-A(K2)
      A(K1)=A(K1)+A(K3)
      A(K3)=A(K1)-A(K3)
      A(K2)=A(K1)+A(K1)
      A(K1)=A(K1)-A(K1)
      A(K2)=A(K2)+A(K3)+1
      A(K3)=A(K2)-A(K3)+1
      C
      T=A((K2)
      A(K2)=A(K1)-T
      A(K1)=A(K1)+T
      V=A((2+1)
      A(K2+1)=A((K+1))-T
      A((K+1))=A((K+1))+T
      C
      T=A(K3)
      A(K3)=A(K1)-T
      A(K1)=A(K1)+T
      V=A((3+1)
      A(K3+1)=A((K1+1))-T
      A((K1+1))=A((K1+1))+T
      C
      T=A(K1)
      A(K1)=A(K1)-T
      A(K1)=A(K1)+T
      V=A((K1+1)
      A(K1+1)=A((K1+1))-T
      A((K1+1))=A((K1+1))+T
      C
      R=-A(K3+1)
      T = A(K3)
      A(K3)=A(K2)-R
      A(K2)=A(K2)+R
      A(K3+1)=A(K2+1))-T
      B0 A(K2+1)=A((K2+1))+T
      B0 IF ((J, LAST) 235.235.62
      B2 J=JDIF +1
      C
      DO FOR J=1
      BLAST= IL +JJ
      DO 85 I = JJ+LAST, IDIF
      KLAST = KL+1
      DO 85 K=1, KLAST, 2
      K1 = K+KBIT
      K2 = K1+KBIT
      K3 = K2+KBIT
      C
      ISN 0061
      ISN 0062
      ISN 0063
      C
      ISN 0064
      ISN 0065
      ISN 0066
      ISN 0067
      ISN 0068
      ISN 0069
      C
      ISN 0070
      ISN 0071
      ISN 0072
      ISN 0073
      ISN 0074
      ISN 0075
      C
      ISN 0076
      ISN 0077
      ISN 0078
      ISN 0079
      ISN 0080
      ISN 0081
      C
      ISN 0082
      ISN 0083
      ISN 0084
      ISN 0085
      ISN 0086
      ISN 0087
      ISN 0088
      ISN 0089
      C
      ISN 0090
      ISN 0091
      ISN 0092
      ISN 0093
      ISN 0094
      ISN 0095
      ISN 0096
      C
      HARM1610
      HARM1620
      HARM1630
      HARM1640
      HARM1650
      HARM1660
      HARM1670
      HARM1680
      HARM1690
      HARM1700
      HARM1710
      HARM1720
      HARM1730
      HARM1740
      HARM1750
      HARM1760
      HARM1770
      HARM1780
      HARM1790
      HARM1800
      HARM1810
      HARM1820
      HARM1830
      HARM1840
      HARM1850
      HARM1860
      HARM1870
      HARM1880
      HARM1890
      HARM1900
      HARM1910
      HARM1920
      HARM1930
      HARM1940
      HARM1950
      HARM1960
      HARM1970
      HARM1980
      HARM1990
      HARM2000
      HARM2010
      HARM2020
      HARM2030
      HARM2040
      HARM2050
      HARM2060
      HARM2070
      HARM2080
      HARM2090
      HARM2100
      HARM2110
      HARM2120
      HARM2130
      HARM2140
  
```

LEVEL 2.30 - (JUNE 78) HARM OS/3.0 JETRIAN M EXTENDED DATE 80.194 1.57.09

C LEAVING M=(1+1)/R0012. M3=(-1+1)/20J12.M2=1.

C A(K1)=A(K1)*A(K2)+1
 C A(K2)=A(K1)-A(K2)+1
 C A(K1)=A(K1)*M+A(K3)*M
 C A(K3)=A(K1)*M-A(K3)*M
 C A(K1)=A(K1)+A(K1)
 C A(K1)=A(K1)-A(K1)
 C A(K2)=A(K2)+A(K3)+1
 C A(K3)=A((2)-A(K3)+1
 C R =-A(K2+1)
 C T = A(K2)
 C A(K2) = A(K)-R
 C A(K1) = A(K)+R
 C A(K2+1)=A(K+1)-T
 C A(K+1)=A(K+1)+T

C A(R)=A(K1)-A(K1+1)
 C A(M) = A((1+1)+1)(K1)
 C R=A(K3)-A(K3+1)
 C T=A(K3)-A(K3+1)
 C A(K3)=A(R-R)/ROOT2
 C A(K3+1)=(A(R1-T)/ROOT2
 C A(K3)=A(WR-R)/ROOT2
 C A(K3)=A(WR-R)/ROOT2
 C ISN 0103
 C ISN 0104
 C ISN 0105
 C ISN 0106
 C ISN 0107
 C ISN 0108
 C ISN 0109
 C ISN 0110
 C ISN 0111
 C ISN 0112
 C ISN 0113
 C ISN 0114
 C ISN 0115
 C ISN 0116
 C ISN 0117
 C ISN 0118
 C ISN 0119
 C ISN 0120
 C ISN 0121
 C ISN 0122
 C ISN 0123
 C ISN 0124
 C ISN 0125

C A(K1)=A(K1)+1
 C A(K2)=A(K+1)-T
 C A(K1+1)=A(K+1)+T
 C R=-A(K3+1)
 C T=A((3))
 C A(K3)=A(K2)-R
 C A(K2)=A(K2)+R
 C A(K3+1)=A(K2+1)+T
 C A(K2+1)=A(K2+1)+T
 C 85 IF(JJ-LAST-1) 235.235.90
 C 90 JJ + JJ0 IF
 C NOW DO THE REMAINING J'S
 C DO 230 J=2, JLAST

C C FETCH M*S
 C C DEF- M=M**INV(J). M2=M**2. M3=M**3
 C C 96 I=INV(J+1)
 C C 98 I=NT-1
 C C M(1)=S(1C)
 C C M(2)=S(1)
 C C 12=2+1
 C C 12C=NT-12
 C C IF((12C)120.110.100
 C C ISN 0126
 C C ISN 0127
 C C ISN 0128
 C C ISN 0129
 C C ISN 0130
 C C ISN 0131
 C C ISN 0132

LEVEL 2. (JUNE 76) HARM OS/350 FORTRAN H EXTENDED DATE 80.19 21.57.09

201 IS IN FIRST QUADRANT
 ISN 0133 C 100 W2(1)=S(12C)
 ISN 0134 W2(2)=S(12C)
 ISN 0135 GO TD 130
 ISN 0136 110 W2(1)=0.
 ISN 0137 W2(2)=1.
 ISN 0138 GO TD 130

C 201 IS IN SECOND QUADRANT
 ISN 0139 C 120 12CC = [2C+N]
 ISN 0140 12C = 12C
 ISN 0141 W2(2)=S(12CC)
 ISN 0142 130 13+12
 ISN 0143 13C-N=13
 ISN 0144 F((13C))160+150+140
 ISN 0145 C 13 IN FIRST QUADRANT
 ISN 0146 C 140 W3(1)=S(13C)
 ISN 0147 W3(2)=S(13C)
 ISN 0148 GO TD 200
 ISN 0149 150 W3(1)=0.
 ISN 0150 W3(2)=1.
 ISN 0151 GO TD 200
 ISN 0152 C 160 13CC=13C.MT
 ISN 0153 C 170 13 IN SECOND QUADRANT
 ISN 0154 ISN 0155 C 170 13C=13C
 ISN 0156 W3(1)=S(13C)
 ISN 0157 GO TD 200
 ISN 0158 180 W3(1)=1.
 ISN 0159 W3(2)=0.
 ISN 0160 GO TD 200

C 301 IS IN THIRD QUADRANT
 ISN 0161 ISN 0162 C 190 13CC=N+13CC
 ISN 0163 W3(1)=S(13CC)
 ISN 0164 W3(2)=S(13CC)
 200 1LAS=1L+JJ
 ISN 0165 00 220 1=JJ.1LAST.1DIF
 ISN 0166 KLAST.KL+1
 ISN 0167 00 220 K=1.KLAST.2
 ISN 0168 K1=K+K3IT
 ISN 0169 K2=K1+K6IT
 ISN 0170 K3=K2+K8IT
 ISN 0171 C
 C 00 TWO STEPS WITH J NOT 0
 A(K1)=A(K1)+A(K2)*W2
 A(K2)=A(K2)+A(K1)*W2
 A(K1)=A(K1)+A(K3)*W3
 C

LEVEL 2.3.3 (JUNE 78) HARM 05/360 JETRAN H EXTENDED DATE 80.195 21.57.09

```

C A(K3)=A(K1)*B(K3)+A(K3)*B3 HARM3230
C A(K1)=A(K1)+A(K1) HARM3240
C A(K1)=A(K1)-A(K1) HARM3250
C A(K2)=A(K2)+A(K3)+I HARM3260
C A(K3)=A(K2)-A(K3)+I HARM3270
C R=A(K2)*B2(1)-A(K2)*B2(2) HARM3280
C T=A(K2)*B2(2)+A(K2)+B2(1) HARM3290
C A(K2)=A(K1)-R HARM3300
C A(K1)=A(K1)+R HARM3310
C A(K2+1)=A(K1)+T HARM3320
C A(K1+1)=A(K1)+T HARM3330
C R=A(K3)*B3(1)-A(K3)*B3(2) HARM3340
C T=A(K3)*B3(2)+A(K3)+B3(1) HARM3350
C A(K1)=A(K1)*B1(1)-A(K1)*B1(2) HARM3360
C A(K1)=A(K1)*B1(2)-A(K1)*B1(1) HARM3370
C A(K3)=A82-R HARM3380
C A(K3+1)=AM1-T HARM3390
C A(K1)=AMR+R HARM3400
C A(K1+1)=AM1+T HARM3410
C T=A(K1) HARM3420
C A(K1)=A(K1)-T HARM3430
C A(K1)=A(K1)+T HARM3440
C T=A(K1+1) HARM3450
C A(K1+1)=A(K1+1)-T HARM3460
C A(K1+1)=A(K1+1)+T HARM3470
C R=A(K3+1)-T HARM3480
C A(K3+1)=AM1-T HARM3490
C A(K1+1)=A(K1+1)+T HARM3500
C T=A(K3+1) HARM3510
C A(K3)=A(K2)-R HARM3520
C A(K2)=A(K2)+R HARM3530
C A(K2+1)=A(K2+1)-T HARM3540
C A(K2+1)=A(K2+1)+T HARM3550
C 220 END 3F 1 AND K LOOPS HARM3560
C
C 230 JJ=JJ0+JJ HARM3580
C END 3F J-LOOP HARM3600
C
C 0199 235 JLAST=JJLAST+3 HARM3610
C 0200 240 CONTINUE HARM3620
C END OF L LOOP HARM3630
C
C 0201 250 CONTINUE HARM3640
C END OF ID LOOP HARM3650
C
C 0202 360 GO3=1 HARM3660
C 0203 360 GO3=0, DM EQ, MT HARM3670
C 0204 360 GO3=1 HARM3680
C
C WE NOW HAVE THE COMPLEX FOURIER SUMS BUT THEIR ADDRESSES ARE
C BIT-REVERSED. THE FOLLOWING RJSINE PUTS THEM IN ORDER
C NTSQ=NTRANS HARM3700
C N3MT=M3-47 HARM3720
C 350 IF(M3MT) 370,360,360 HARM3730
C
C 360 GO3=1 HARM3740
C
C 0205 360 GO3=1 HARM3750
C

```

LEVEL 2.3.u (JUNE 78)	HARM	05/300	ATTAN H EXTENDED	DATE 00.199. • 57.00
1SN 0206	N3VNT=N3/NT			HARM3770
1SN 0207	MINN3=NT			HARM3780
1SN 0208	GO TO 360	C		HARM3790
1SN 0209	C 43 LESS THAN MT	370	I GO3=2	HARM3810
1SN 0210	N3VNT=1		NTVN3=NT/N3	HARM3830
1SN 0211	MINN3=N3			HARM3840
1SN 0212	380 JJD3 = NT50/N3			HARM3850
1SN 0213	42M=42-MT			HARM3860
1SN 0214	450 IF (N2NT)470.460.460	C		HARM3870
1SN 0215	M2 GR. OR EQ. MT	460	I GO2=1	HARM3880
1SN 0216	N2VNT=N2/NT		N2VNT=N2/NT	HARM3890
1SN 0217	MINN3=NT		NTVN2=NT/N2	HARM3910
1SN 0218	GO TO 480	C		HARM3920
1SN 0219	C M2 LESS THAN MT	470	I GO2 = 2	HARM3930
1SN 0220	N3VNT=1		N3VNT=N2	HARM3940
1SN 0221	MINN3=N2		NTVN2=NT/N2	HARM3950
1SN 0222	480 JJD2=NT50/N2			HARM3960
1SN 0223	MIN1=NT	C	MIN1=NT	HARM3970
1SN 0224	550 IF(M14)570.560.560			HARM3980
1SN 0225	C M1 GR. OR EQ. MT	560	I GO1=2	HARM3990
1SN 0226	N1VNT=N1/NT		N1VNT=N1/NT	HARM4000
1SN 0227	MINN1=NT		NTVN1=NT/N1	HARM4010
1SN 0228	GO TO 580	C	MIN1=NT	HARM4020
1SN 0229	C M1 LESS THAN MT	570	I GO1=2	HARM4030
1SN 0230	N1VNT=N1/NT		N1VNT=N1/NT	HARM4040
1SN 0231	MINN1=NT		NTVN1=NT/N1	HARM4050
1SN 0232	GO TO 580	C	MINN1=NT	HARM4070
1SN 0233	C M1 LESS THAN MT	580	I GO1=2	HARM4080
1SN 0234	N1VNT=N1/NT		N1VNT=N1/NT	HARM4090
1SN 0235	600 JJD1=NT50/41	C	MINN1=NT	HARM4100
1SN 0236	J=1			HACK4110
1SN 0237	DD 890 JPP3=1.N3VNT			HARM4120
1SN 0238	I PP3=1.NV(JJ3)			HARM4130
1SN 0239	DD 870 J23=1.MINN3			HARM4140
1SN 0240	GO T2 (610.620) I GO3			HARM4150
1SN 0241	610 I P3=1.NV(JP3)*N3VNT			HARM4160
1SN 0242	GO T2 610			HARM4170
1SN 0243	620 I P3=1.NV(JP3)/NTVN3			HARM4180
1SN 0244	630 I 3=(I P3+IP3)*N2			HARM4190
1SN 0245	700 J J2=1			HARM4200
1SN 0246	DD 870 J23=1.N2VNT			HARM4210
1SN 0247	I PP2=1.NV(JJ2)+13			HARM4220
1SN 0248	DD 850 J2=1.M1NN2			HARM4230
1SN 0249				HARM4240


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LEVEL 2.3.~ ( JUNE 79)      HARM
                                OS/350      ARTRAN H EXTENDED      DATE 80.199.      .57.09

ISN 0286      S(JDIF)=SINT(META)
ISN 0289      DO 950 L=2,M
ISN 0290      THE TA=THE TA/F.
ISN 0291      JSTEP=2*JSTEP
ISN 0292      JDIF=JDIF
ISN 0293      JDIF=JSTEP/2
ISN 0294      SJDF1=SIN(META)
ISN 0295      JC1=N-JDIF
ISN 0296      SJCI=CCS(META)
ISN 0297      JLAST=NT-JSTEP
ISN 0298      IF (JLAST-N-JSTEP) 950,920,920
ISN 0299      DO 940 J=JSTEP,JLAST,JSTEP
ISN 0300      JC=N-
ISN 0301      JD=J+JDIF
ISN 0302      SJDI=S(J)*SJCI+S(JDIF)*SJJC
ISN 0303      950 CONTINUE
ISN 0304      C      SET J> INV(J) TABLE
ISN 0305      C      MTLEXP=2**(NT-L). FOR L=1
ISN 0306      L1EXP=1
ISN 0307      L1EXP=2**((L-1). FOR L=1
ISN 0308      INV(L)=0
ISN 0309      DO 940 L=1,M
ISN 0310      INV(M1EXP+1)
ISN 0311      DO 970 J=2,L1EXP
ISN 0312      INV(J+1)EXP
ISN 0313      INV(J)=INV(J)*MTLEXP
ISN 0314      MTLEXP=MTLEXP/2
ISN 0315      L1EXP=L1EXP/2
ISN 0316      980 IF (L1EXP) 12,955,12
ISN 0317      END

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTOBLB(NONE)
OPTIONS IN EFFECT: SOURCE EBCDIC NOLIST NODECK OBJECT 'V4AP' NOFORMAT GOSTMT NOXREF ALC NOANSF NUTERM 'ION FL'
*STATISTICS* SOURCE STATEMENTS = 314. PROGRAM SIZE = 5206. SUBPROGRAM NAME = HARM
*STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPIILATION *****

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56K BYTES OF CORE NOT USED


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LEVEL 2-3-0 (JUNE 76)      GREYSC      OS/360      RTAN H EXTENDED      DATE 80-199.  57-13

1SN 0044      DO 400 I=1,MIN,1,MAX,1,DEL
1SN 0045      DO 250 I=1,3
1SN 0046      DO 250 J=2,NLASTL
1SN 0047      250 LINE(J,I1)=IBLANK
1SN 0048      LEV=1
1SN 0049      J=2
1SN 0050      DO 300 JJ=JMIN,JMAX,JDEL
1SN 0051      J=J+1
1SN 0052      L=(AMAT((J,J))-ANN)/(AMX-ANN)*FLDAT(INLEVEL)+1.
1SN 0053      L=MAXD(I1,L)
1SN 0054      L=MIND(INLEVEL,L)
1SN 0055      LEVMAXD(LEVEL,L1)
1SN 0056      LEVNONLE(LEVEL,L1)
1SN 0057      DO 300 K=1,LEVELNOW
1SN 0058      LINE(I,J,K)=CHARS(L,K)
1SN 0059      300 CONTINUE
C
C      FIND LAST PRINT POSITION
C
C      DO 400 KL=1,LEVEL
C      DO 350 K=1,LASTL
C      IF(LINE(K,KL)=NF.16BLANK) G2 YJ 375
C      CONTINUE
C      IF(ILEX.GT.0.AND.KL.EQ.LEV) WRITE(FILE,1050) (LINE(I1,KL),
C      375
C      350
C      375
C      111=1,KK)
C      IF(ILEX.LT.0.AND.KL.EQ.1) WRITE(FILE,1050) (LINE(I1,KL),
C      111=1,KK)
C      IF(ILEX.GT.0.AND.KL.NE.1) WRITE(FILE,1060) (LINE(I1,KL),
C      111=1,KK)
C      IF(ILEX.LT.0.AND.KL.NE.1) WRITE(FILE,1060) (LINE(I1,KL),
C      111=1,KK)
C      400 CONTINUE
C
C      WRITE LAST LINE (BORDER)
C
C      WRITE(FILE,1070) (IBORDR,I=1,LASTL)
C
C      PRINT TRAILING SCALE INFORMATION
C
C      IF(INCPL.EQ.6) WRITE(FILE,1030) (TITLE(I1),I=1,NWORDS)
C      IF(INCPL.EQ.4) WRITE(FILE,1001) (TITLE(I1),I=1,NWORDS)
C      IF(INCPL.EQ.10) WRITE(FILE,1022) (TITLE(I1),I=1,NWORDS)
C      IF(INCPL.EQ.5) WRITE(FILE,1033) (TITLE(I1),I=1,NWORDS)
C      WRITE(FILE,1010)
C      DELTA=(AMX-ANN)/FLOAT(LEVEL)
C      DO 200 J=1,LEVEL
C      XMIN=FLDAT(I1,FLOAT(LEVEL)*(I1-1)*(A4X-ANN)+ANN
C      XMAX=XMIN*ELLA
C      LFLEVEL=11
C      DO 175 JE:LEV
C      175 IF(JE:LEV.GT.0.AND.J.EQ.LEV) WRITE(FILE,1020) ICHARS((I,J),XMIN,
C      1XMAX

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LEVEL 2.3. - (JUNE 76) GREYSC OS/360 JERRIAN H EXTENDED DATE 60-194 1-57-13

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ISN 0094      IF(IIFI-EX.LT.0.AND.J.J.EQ.1) WRITE(1,FILE,1020) ICHARS(1,J),X4IN,
ISN 0096      I      XMAX
ISN 0098      IF(IIFI.LEX.GT.0.AND.J.J.NE.0) WRITE(1,FILE,1030) ICHARS(1,J)
ISN 0100      175      CONTINUE
ISN 0101      200      CONTINUE
ISN 0102      RETURN
ISN 0103      1000      FORMAT(10-20A6)
ISN 0104      1001      FORMAT(10-20A4)
ISN 0105      1002      FORMAT(10-12A10)
ISN 0106      1003      FORMAT(10-24A5)
ISN 0107      1010      FORMAT(10-32HGREY-SCALE CHARACTERS AND RANGFS/1X)
ISN 0108      1020      FORMAT(5-6A1.5X,E15.6,5X,E15.6)
ISN 0109      1030      FORMAT(11-4,4A,4I)
ISN 0110      1039      FORMAT(11-4)
ISN 0111      1040      FORMAT(11/1X,20A6)
ISN 0112      1041      FORMAT(11/1X,20A4)
ISN 0113      1042      FORMAT(11/1X,12A10)
ISN 0114      1043      FORMAT(11/1X,24A5)
ISN 0115      1045      FORMAT(11-30,132A1)
ISN 0116      1050      FORMAT(11X,132A1)
ISN 0117      1060      FORMAT(1H,132A1)
ISN 0118      1070      FORMAT(1X,132A1)
ISN 0119      END

OPTIONS IN EFFECT: NAME (MAIN) JPTINIZ(2) LINECOUNT(55) SIZE(MAX) AUTODR(NONE)

OPTIONS IN EFFECT: SOURCE EBCDIC NULIST NODECK OBJECT NUMAP NOFORMAT GOSTMT NOKHEZ ALC NOANSF NOTERN FRA S1A
*STATISTICS*      SOURCE STATEMENTS = 118. PROGRAM SIZE = 4040. SUBPROGRAM NAME =GREYSC
*STATISTICS*      NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****

80K BYTES OF CORE NOT USED

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LEVEL 2.3.0 (JUNE 78)

STATISTICS NO DIAGNOSTICS THIS STEP

05/360 FORTRESS H EXTENDED DATE 80.194, 21-57.15

360 LOADER
OPTIONS USED - PRINT, MAP, NOLET, CALL, RES, NOTERM, SIZE=33222, NAME=GD

NAME	TYPE	ADDR	NAME	TYPE	ADDR	NAME	TYPE	ADDR	NAME	TYPE	ADDR
MAIN	SD	116010	PEAK	SD	117406	312E	SD	1176F8	ARRAYS	SD	117410
PRINTL	SD	137468	SD	137490	3REYAC	SD	138EE8	PRINTER	SD	13A1D0	
FLXPI#	LR	13A488	SD	13A508	1HUECOVH	LR	13A604	1ACUME	*	LR	13A604
INTSMTC1#	LR	13B3C0	SD	13B488	FIG12#	LR	13B9FC	1HOCOM442#	SD	13B3B0	
1HOCVTH#	SD	13C548	SD	13C548	420N#	LR	13CSF2	FCVLOUTP	SD	13C682	
FCVLOUTP	LR	13CBBA	FCVDOJTP#	LR	13CCAA	FCVJU12#	LR	13CF0C	1HOLF11#		
ARITH#	LR	13CF90	425EWTC1#	LR	13D524	1HUEF115#	LR	13D790	FLDC3#		
1HOFIUS2#	SD	13F920	1HORZON1#	SD	13EF68	1UCD119#	SD	13F268	FOCUNO#		
1HDERM#	SD	13F730	ERRON1#	LR	13F730	1HJ23RE	SD	13F758	FTEN#		
1HOTRCM#	SD	13FEF0	140TRC1#	LR	13FEF0	3RN14A	SD	1401A0	1HQUAT#L		
1HDNAMEL#	SD	140C10	FRDY#	LR	140C10	EWRN4#	SD	1418A8	FRXPI#		
1HOSURT#	SD	141A26	SORF#	LR	141A26	1HS2JRF	*	SD	1413U0	COS#	
1HOSACOS#	LR	141B90	SIN#	LR	141BB2	1HS3SIN#	SD	141098	THSFX#		
EXP#	LR	141D9A	LC42	CW	141F48						
TOTAL LENGTH 5C736											
ENTRY ADDRESS 116010											

P-10 - 1710x

REFND X

REFNDX

GREY-SCALE CHARACTERS AND RANGES

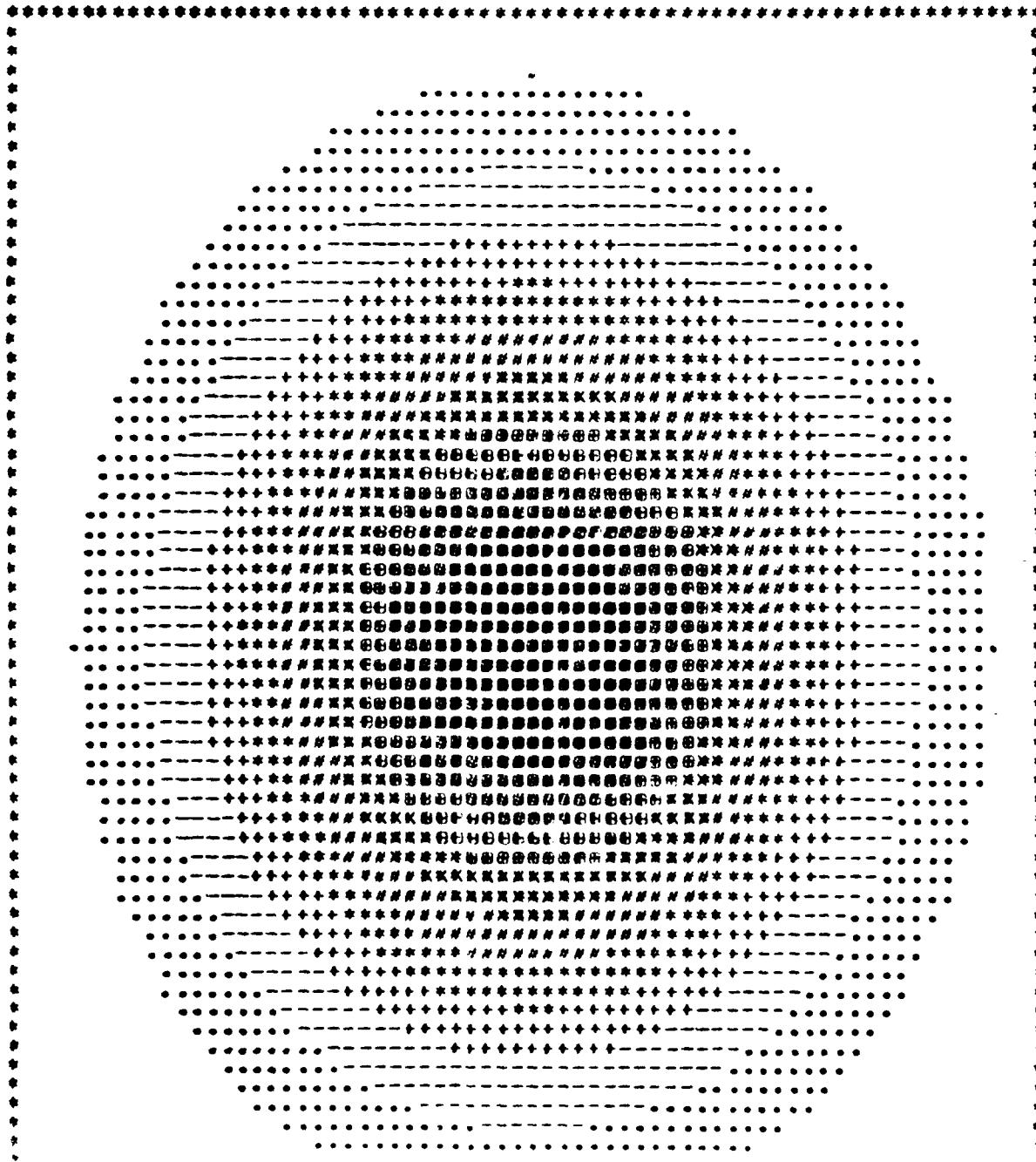
•	0.145928E+01	0.146135E+01
-	0.146135E+01	0.146342E+01
+	0.146342E+01	0.146549E+01
*	0.146550E+01	0.146757E+01
/	0.146757E+01	0.146964E+01
#	0.146964E+01	0.147171E+01
K	0.147171E+01	0.147378E+01
█	0.147378E+01	0.147585E+01
█	0.147585E+01	0.147793E+01
█	0.147793E+01	0.148000E+01

THIS IS STEP 0

SUML = 0.4860770E+00 SUMR = 0.5139230E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.99481 19.99631 MICRONS

IRRADIANT



IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

•	0.597604E-02	0.105378E+00
-	0.105378E+00	0.204781E+00
+	0.204781E+00	0.304183E+00
*	0.304183E+00	0.403585E+00
#	0.403586E+00	0.502988E+00
*	0.502988E+00	0.602390E+00
®	0.602390E+00	0.701793E+00
■	0.701793E+00	0.801195E+00
■	0.801195E+00	0.900598E+00
■	0.900598E+00	0.100000E+01

THIS IS STEP 1

SUML = 0.4859293E+00 SUMR = 0.5140707E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.50356 19.50494 MICRONS

A decorative border consisting of a grid of stars and dots. The border is approximately 100 units wide and 10 units high. It features a central vertical column of stars, a horizontal row of stars at the top, and a decorative pattern of stars and dots along the sides and bottom.

ERRAD CAN

GREY-SCALE CHARACTERS AND RANGES

0.585011E-02	0.105238E+00
0.105238E+00	0.204626E+00
0.204626E+00	0.304013E+00
0.304013E+00	0.403401F+00
0.403401E+00	0.502789E+00
0.502789E+00	0.602177E+00
0.602177E+00	0.701565E+00
0.701565E+00	0.800952E+00
0.800952E+00	0.900340E+00
0.900340E+00	0.999728E+00

THIS IS STEP 2

SUML = 0.4951605E+00 SUMR = 0.5148394E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.72375 18.72428 MICRONS

P-175c

IRRADIATION

IRRADIANT

GREY-SCALE CHARACTERS AND RANGES

0	0.323949E-02	0.211144E+00
.	0.211144E+00	0.419048E+00
-	0.419048E+00	0.626953E+00
*	0.626953E+00	0.834857E+00
+	0.834857E+00	0.104276E+01
#	0.104276E+01	0.125066E+01
%	0.125067E+01	0.145857E+01
•	0.145857E+01	0.166647E+01
•	0.166647E+01	0.187438E+01
■	0.187438E+01	0.208228E+01

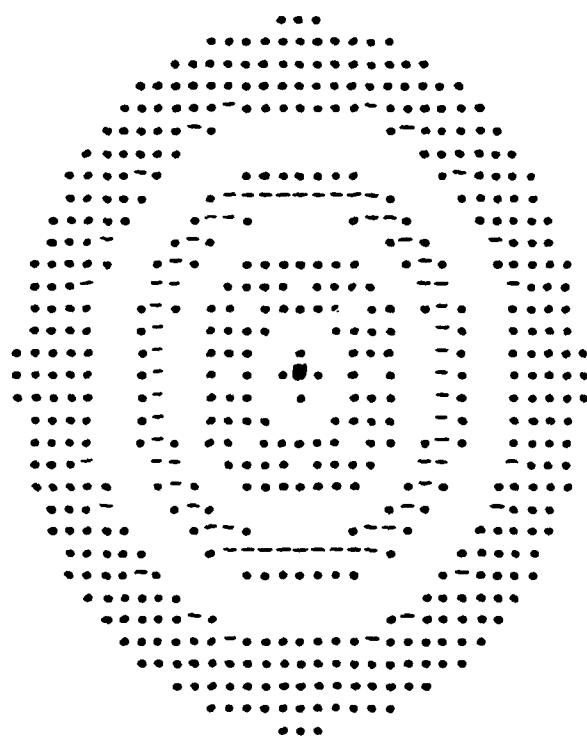
THIS IS STEP 3

SUML = 0.4817133E+00 SUMR = 0.5182866E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.07376 18.07498 MICRONS

P-162

IRRADIANT



A-36

IRRADIANCE

GREY-SCALE CHARACTERS AND RANGES

0.759699E-05	0.842067E+00
0.842067E+00	0.168413E+01
0.168412E+01	0.252618E+01
0.252618E+01	0.336824E+01
0.336824E+01	0.421030E+01
0.421030E+01	0.505236E+01
0.505236E+01	0.589442E+01
0.589442E+01	0.673648E+01
0.673648E+01	0.757854E+01
0.757854E+01	0.842060E+01

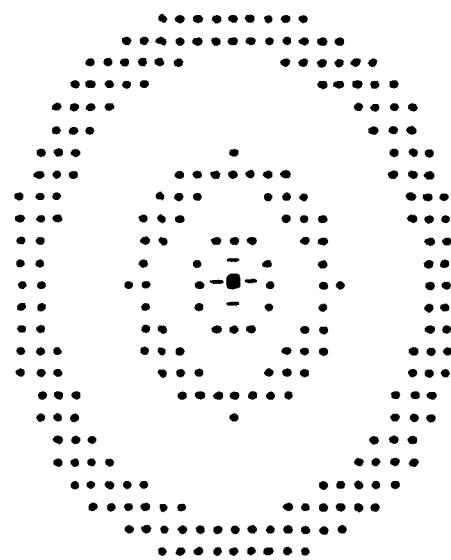
THIS IS STEP 4

SUML = 0.4754750E+00 SUMR = 0.5245249E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 17.64449 17.64572 MICRONS

7-10

IRRADIAN



IRRADIATION

GREY-SCALE CHARACTERS AND RANGES

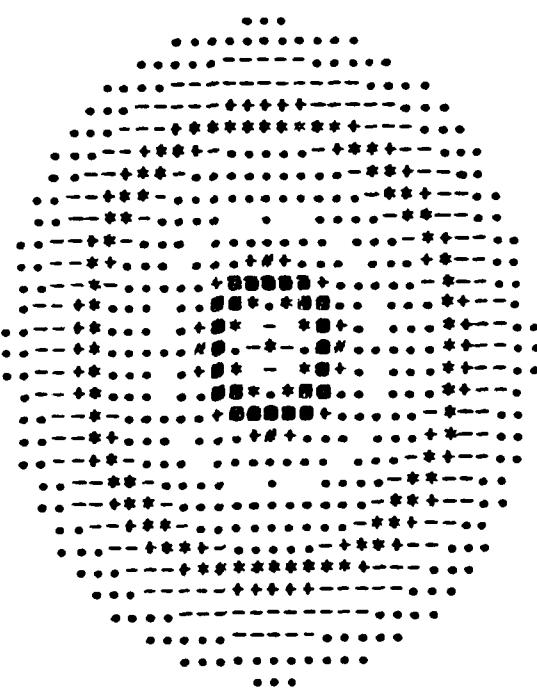
0.247009E-03	0.153054E+01
0.153054E+01	0.306083E+01
0.306083E+01	0.459112E+01
0.459112E+01	0.612142E+01
0.612142E+01	0.765171E+01
0.765171E+01	0.918200E+01
0.918200E+01	0.107123E+02
0.107123E+02	0.122426E+02
0.122426E+02	0.137729E+02
0.137729E+02	0.153032E+02

THIS IS STEP 5

SUML = 0.4803946E+00 SUMR = 0.5196053E+00

SEAM WAIST IN X AND Y DIRECTIONS IS 17.43616 17.43765 MICRONS

ERRADIAN



IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

•	0.348578E-05	0.546719E+00
-	0.546719E+00	0.109343E+01
+	0.109343E+01	0.154015E+01
*	0.164015E+01	0.218686E+01
#	0.218686E+01	0.273358E+01
%	0.273358E+01	0.328029E+01
€	0.328029E+01	0.382701E+01
¤	0.382701E+01	0.437372E+01
¤	0.437373E+01	0.492044E+01
¤	0.492044E+01	0.546716E+01

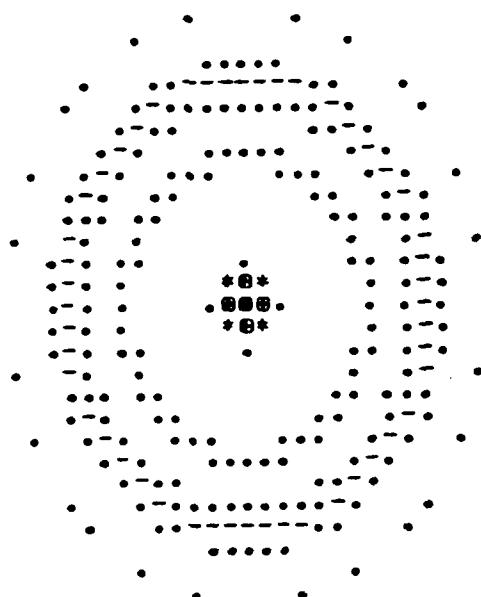
THIS IS STEP 6

SUML = 0.4704280E+00 SUMR = 0.5295719E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 17.47115 17.47253 MICRONS

P-12-1

IRRADIAN



ERRADIAN

GREY-SCALE CHARACTERS AND RANGES

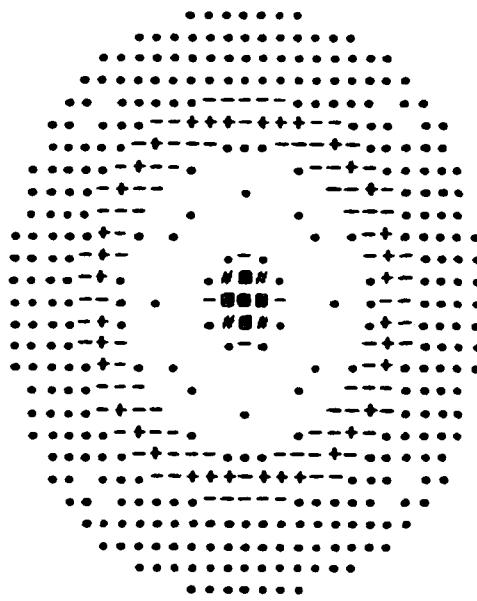
0	0.341598E-04	0.146743E+01
.	0.146743E+01	0.293483E+01
-	0.293483E+01	0.440223E+01
+	0.440223E+01	0.586963E+01
*	0.586963E+01	0.733703E+01
#	0.733703E+01	0.880443E+01
R	0.880443E+01	0.102718E+02
•	0.102718E+02	0.117392E+02
■	0.117392E+02	0.132066E+02
■	0.132066E+02	0.146740E+02

THIS IS STEP 7

SUML = 0.4736281E+00 SUMR = 0.5263718E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 17.75745 17.75934 MICRONS

IRRADIANT



IRRADIATION

GREY-SCALE CHARACTERS AND RANGES

0.349578E-03	0.996981E+00
0.996981E+00	0.149361E+01
0.199361E+01	0.299024E+01
0.299024E+01	0.398687E+01
0.398688E+01	0.498351E+01
0.498351E+01	0.598014E+01
0.598014E+01	0.697677E+01
0.697677E+01	0.797340E+01
0.797340E+01	0.897003E+01
0.897003E+01	0.996667E+01

THIS IS STEP 8

SUML = 0.4800833E+00 SUMR = 0.5199167E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.20370 18.20607 MICRONS

IRRADIATION

IRRADIANT

GREY-SCALE CHARACTERS AND RANGES

0.213262E-03	0.358114E+00
0.358113E+00	0.716014E+00
0.716014E+00	0.107391E+01
0.107391E+01	0.143181E+01
0.143181E+01	0.178971E+01
0.178971E+01	0.214761E+01
0.214761E+01	0.250551E+01
0.250551E+01	0.286341E+01
0.286341E+01	0.322131E+01
0.322132E+01	0.357922E+01

THIS IS STEP 9

SUML = 0.4812048E+00 SUMR = 0.5187951E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.54669 18.54912 MICRONS

IRRADIATION

IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

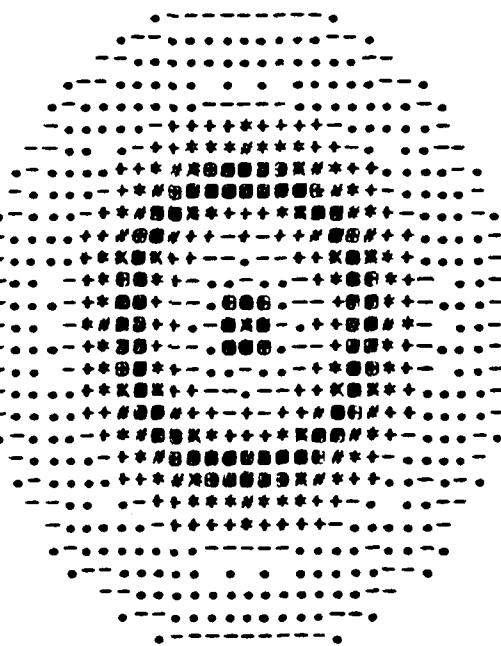
•	0.753595E-04	0.379830E+00
-	0.379830E+00	0.759584E+00
+	0.759584E+00	0.113934E+01
*	0.113934E+01	0.151909E+01
#	0.151909E+01	0.189885E+01
%	0.189885E+01	0.227860E+01
R	0.227860E+01	0.265836E+01
■	0.265836E+01	0.303811E+01
●	0.303811E+01	0.341786E+01
○	0.341787E+01	0.379762E+01

THIS IS STEP 10

SUM_ = 0.4789906E+00 SUMR = 0.5210093E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.63527 18.63782 MICRONS

IRRADIANT



* * * * *
IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.103213E-03	0.410842E+00
-	0.410842E+00	0.821581E+00
*	0.821582E+00	0.123232E+01
+	0.123232E+01	0.164306E+01
*	0.164306E+01	0.205380E+01
/	0.205380E+01	0.246454E+01
R	0.246454E+01	0.287528E+01
0	0.287528E+01	0.328602E+01
0	0.328602E+01	0.369676E+01
0	0.369676E+01	0.410749E+01

THIS IS STEP 11

SUML = 0.4775044E+00 SUMR = 0.5224955E+00

BEAN WAIST IN X AND Y DIRECTIONS IS 18.77403 18.77573 MICRONS

IRRADIATION

IRRADIATION

GREY-SCALE CHARACTERS AND RANGES

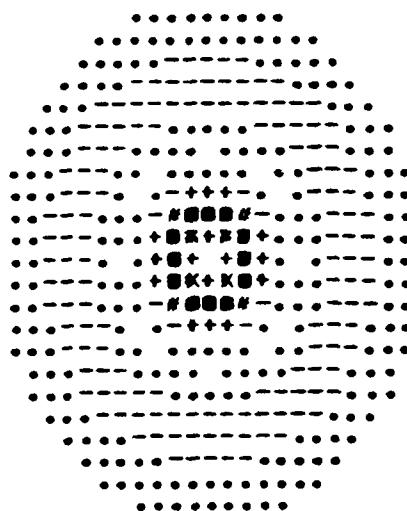
0.373785E-04	0.594157E+00
0.594157E+00	0.118828E+01
0.118828E+01	0.178239E+01
0.178239E+01	0.237651E+01
0.237651E+01	0.297063E+01
0.297063E+01	0.356475E+01
0.356475E+01	0.415887E+01
0.415887E+01	0.475299E+01
0.475299E+01	0.534711E+01
0.534711E+01	0.594123E+01

THIS IS STEP 12

SUML = 0.4750217E+00 SUMR = 0.5249783E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.92630 18.92923 MICRONS

IRRADIANT



IRRADIANT

GREY-SCALE CHARACTERS AND RANGES

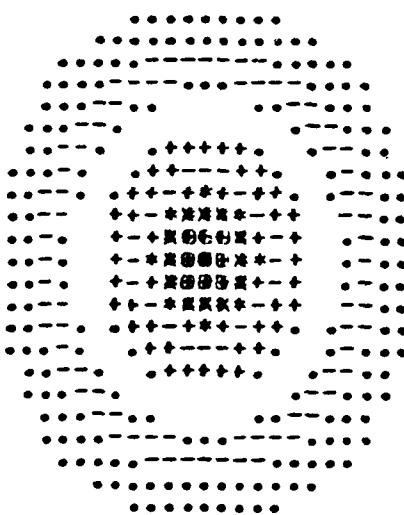
	0.126287E-04	0.898516E+00
.	0.898516E+00	0.179702E+01
-	0.179702E+01	0.269552E+01
◆	0.269552E+01	0.359403E+01
*	0.359403E+01	0.449253E+01
#	0.449253E+01	0.539103E+01
×	0.539103E+01	0.628954E+01
◎	0.628954E+01	0.718804E+01
■	0.718804E+01	0.808654E+01
■	0.808654E+01	0.898505E+01

THIS IS STEP 13

SUML = 0.4702227E+00 SUMR = 0.5297772E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.73090 18.73331 MICRONS

IRRADIAN



IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

0.703742E-04	0.848021E+00
0.848021E+00	0.169597E+01
0.169597E+01	0.254392E+01
0.254392E+01	0.339187E+01
0.339187E+01	0.423982E+01
0.423982E+01	0.508778E+01
0.508778E+01	0.593573E+01
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0.678368E+01	0.763163E+01
0.763163E+01	0.847958E+01

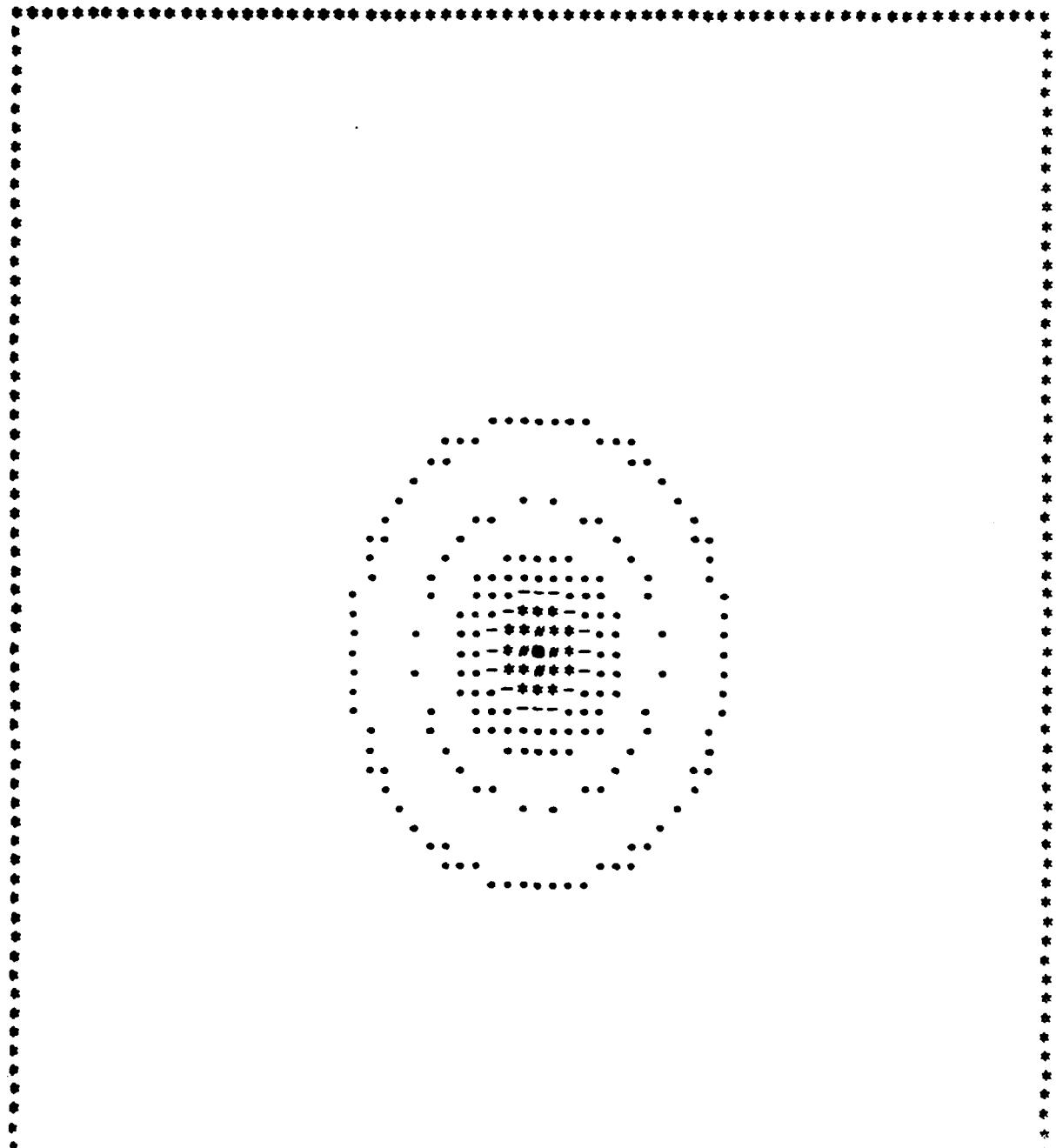
THIS IS STEP 14

SUML = 0.4651438E+00 SUMR = 0.5348561E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.77950 18.78253 MICRONS

1000

ERRADIAN



ERRADIAN

GREY-SCALE CHARACTERS AND RANGES

•	0.244114E-03	0.155652E+01
-	0.155652E+01	0.311279E+01
+	0.311279E+01	0.466906E+01
*	0.466906E+01	0.622533E+01
#	0.622534E+01	0.778161E+01
%	0.778161E+01	0.933788E+01
■	0.933788E+01	0.108942E+02
■	0.108942E+02	0.124504E+02
■	0.124504E+02	0.140067E+02
■	0.140067E+02	0.155630E+02

THIS IS STEP 15

SUML = 0.4690364E+00 SUMR = 0.5309635E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.87981 18.88269 MICRONS

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THEORETICAL ANALYSIS OF MULTIMODE FIBER STRUCTURES. (U)

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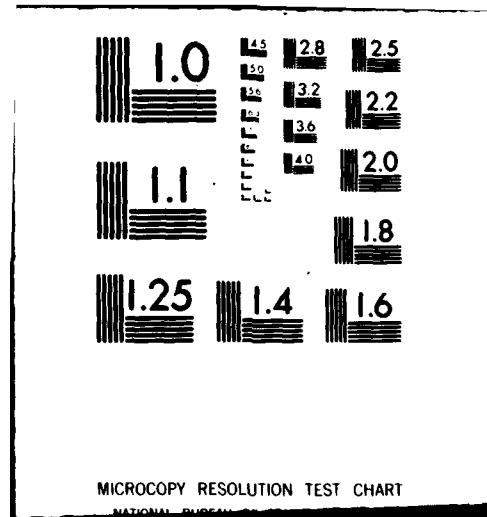
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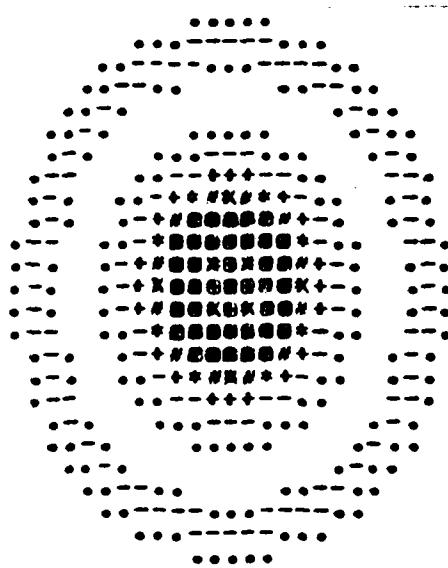


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IRRADIAN



IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

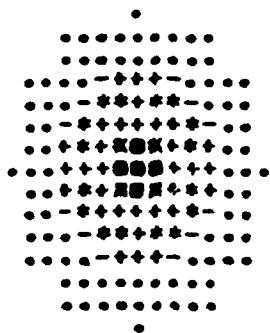
0.246905E-03	0.697666E+00
0.697666E+00	0.139509E+01
0.139508E+01	0.209250E+01
0.209250E+01	0.278992E+01
0.278992E+01	0.348734E+01
0.348734E+01	0.418476E+01
0.418476E+01	0.488218E+01
0.488218E+01	0.557960E+01
0.557960E+01	0.627702E+01
0.627702E+01	0.697444E+01

THIS IS STEP 16

SUML = 0.4632549E+00 SUMR = 0.5367450E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.10669 19.10956 MICRONS

IRRADIANT



ERRADIAN

BREV-SCALE CHARACTERS AND RANGES

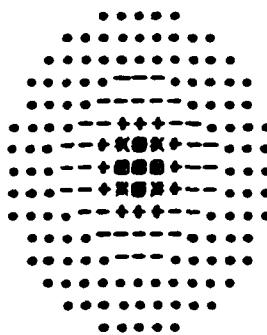
0.911394E-03	0.131676E+01
0.131676E+01	0.263261E+01
0.263261E+01	0.394845E+01
0.394845E+01	0.526430E+01
0.526430E+01	0.658015E+01
0.658015E+01	0.789600E+01
0.789600E+01	0.921185E+01
0.921185E+01	0.105277E+02
0.105277E+02	0.118435E+02
0.118435E+02	0.131594E+02

THIS IS STEP 17

SUML = 0.4618073E+00 SUMR = 0.5381926E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.16985 19.17290 MICRONS

IRRADIANT



IRRADIANT

GREY-SCALE CHARACTERS AND RANGES

0.192583E-04	0.159542E+01
0.159542E+01	0.319081E+01
0.319081E+01	0.478621E+01
0.478621E+01	0.638161E+01
0.638161E+01	0.797700E+01
0.797700E+01	0.957240E+01
0.957240E+01	0.111678E+02
0.111678E+02	0.127632E+02
0.127632E+02	0.143586E+02
0.143586E+02	0.159540E+02

THIS IS STEP 18

SUML = 0.4730968E+00 SUMR = 0.5269032E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.34006 19.34309 MICRONS

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